



Search for direct production of electroweakinos in final states with one lepton, missing transverse momentum and a Higgs boson decaying into two b -jets in pp collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector

ATLAS Collaboration*

CERN, 1211 Geneva 23, Switzerland

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Abstract The results of a search for electroweakino pair production $pp \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$ in which the chargino ($\tilde{\chi}_1^\pm$) decays into a W boson and the lightest neutralino ($\tilde{\chi}_1^0$), while the heavier neutralino ($\tilde{\chi}_2^0$) decays into the Standard Model 125 GeV Higgs boson and a second $\tilde{\chi}_1^0$ are presented. The signal selection requires a pair of b -tagged jets consistent with those from a Higgs boson decay, and either an electron or a muon from the W boson decay, together with missing transverse momentum from the corresponding neutrino and the stable neutralinos. The analysis is based on data corresponding to 139 fb^{-1} of $\sqrt{s} = 13$ TeV pp collisions provided by the Large Hadron Collider and recorded by the ATLAS detector. No statistically significant evidence of an excess of events above the Standard Model expectation is found. Limits are set on the direct production of the electroweakinos in simplified models, assuming pure wino cross-sections. Masses of $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ up to 740 GeV are excluded at 95% confidence level for a massless $\tilde{\chi}_1^0$.

1 Introduction

The Standard Model (SM) is a remarkably successful theory, yet it is clear that this theory is not a complete description of nature. The discovery in 2012 of the SM Higgs boson [1–4], by the ATLAS and CMS collaborations, confirmed the mechanism of the electroweak symmetry breaking and highlighted the hierarchy problem [5–8]. Supersymmetry (SUSY) [9–14], a theoretical extension to the SM, resolves the hierarchy problem by introducing a new fermion (boson) supersymmetric partner for each boson (fermion) in the SM. In SUSY models that conserve R -parity [15], the SUSY particles are produced in pairs. Furthermore, the lightest supersymmetric particle (LSP) is stable and weakly interacting, thus constituting a viable dark-matter candidate [16, 17].

In SUSY scenarios the partners of the SM Higgs boson (h) and the gauge bosons, known as the higgsinos, winos (partners of the $SU(2)_L$ gauge fields), and bino (partner of the $U(1)$ gauge field) are collectively referred to as electroweakinos. Charginos $\tilde{\chi}_i^\pm$ ($i = 1, 2$) and neutralinos $\tilde{\chi}_j^0$ ($j = 1, 2, 3, 4$) are the electroweakino mass eigenstates which are linear superpositions of higgsinos, winos, and bino. For the models considered in this paper, the lightest neutralino ($\tilde{\chi}_1^0$) is a bino-like LSP. The lightest chargino ($\tilde{\chi}_1^\pm$) and next-to-lightest neutralino ($\tilde{\chi}_2^0$) are wino-like and nearly mass degenerate.

Naturalness considerations [18, 19] suggest that the lightest of the electroweakinos have masses near the electroweak scale. In scenarios where the strongly produced SUSY particles are heavier than a few TeV, the direct production of electroweakinos may be the dominant SUSY production mechanism at the Large Hadron Collider (LHC). The lightest chargino and next-to-lightest neutralino can decay via $\tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^0$ and $\tilde{\chi}_2^0 \rightarrow h/Z\tilde{\chi}_1^0$ respectively [20–22] in scenarios where the lepton superpartners are heavier than the electroweakinos. In this case the decay via the Higgs boson is dominant for many choices of SUSY parameters, as long as $m(\tilde{\chi}_2^0) - m(\tilde{\chi}_1^0) > m(h)$. Scenarios with light electroweakinos also provide a possible explanation for the discrepancy between the muon anomalous magnetic moment $g - 2$ measurement and the SM predictions [23, 24].

This paper presents a search for direct production of electroweakinos in proton–proton (pp) collisions produced at the LHC at $\sqrt{s} = 13$ TeV. This analysis is designed to be sensitive to direct production of a chargino and a neutralino that promptly decay as $\tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^0$ and $\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0$. The search targets a W boson which decays into an electron or muon (and corresponding neutrino) and a Higgs boson which decays into a pair of b -quarks, as shown in Fig. 1. The signature consists of exactly one light lepton (e or μ), two jets originating from the fragmentation of b -quarks, and missing

* e-mail: atlas.publications@cern.ch

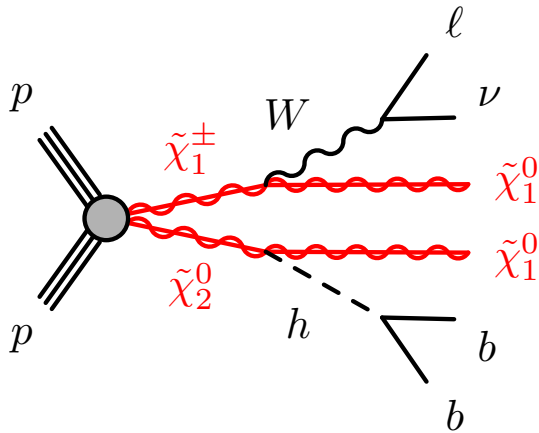


Fig. 1 A diagram illustrating the signal scenario considered for the production of a chargino and a next-to-lightest neutralino

transverse momentum (p_T^{miss}) from neutralinos and neutrinos. A set of simplified SUSY models is used to optimise the search and interpret the results. The branching ratios of $\tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^0$ and $\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0$ are assumed to be 100%. The branching ratio of $h \rightarrow b\bar{b}$ is taken to be 58.3% as expected for the SM Higgs boson.

Previous searches for charginos and neutralinos at the LHC targeting decays via the Higgs boson have been reported by the ATLAS [25] and CMS [26] collaborations. Because of increased integrated luminosity and an improved two-dimensional fit procedure, the search presented here significantly extends the SUSY parameter space sensitivity beyond that of the previously published 13 TeV ATLAS search [25] for the same final state.

2 ATLAS detector

The ATLAS detector [27] is a multipurpose particle detector with a nearly 4π coverage in solid angle.¹ It consists of an inner tracking detector surrounded by a thin superconducting solenoid providing a 2 T axial magnetic field, electromagnetic and hadron calorimeters, and a muon spectrometer. The inner tracking detector covers the pseudorapidity range $|\eta| < 2.5$. It consists of silicon pixel, sili-

¹ ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point in the centre of the detector. The positive x -axis is defined by the direction from the interaction point to the centre of the LHC ring, with the positive y -axis pointing upwards, while the beam direction defines the z -axis. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the z -axis. The pseudorapidity η is defined in terms of the polar angle θ by $\eta = -\ln \tan(\theta/2)$. Rapidity is defined as $y = 0.5 \ln[(E + p_z)/(E - p_z)]$ where E denotes the energy and p_z is the component of the momentum along the beam direction. The angular distance ΔR is defined as $\sqrt{(\Delta y)^2 + (\Delta \phi)^2}$.

con microstrip, and transition radiation tracking detectors. A new inner pixel layer, the insertable B-layer [28, 29], was added at a mean radius of 3.3 cm before the start of 2015 data taking period, improving the identification of b -jets. Lead/liquid-argon (LAr) sampling calorimeters provide electromagnetic (EM) energy measurements with high granularity. A steel/scintillator-tile hadron calorimeter covers the central pseudorapidity range ($|\eta| < 1.7$). The endcap and forward regions are instrumented with LAr calorimeters for EM and hadronic energy measurements up to $|\eta| = 4.9$. The muon spectrometer surrounds the calorimeters and is based on three large air-core toroidal superconducting magnets with eight coils each. The field integral of the toroids ranges between 2.0 and 6.0 Tm across most of the detector. The muon spectrometer includes a system of precision tracking chambers and fast detectors for triggering. A two-level trigger system [30] is used to select events. The first-level trigger is implemented in hardware and uses a subset of the detector information to keep the accepted rate below 100 kHz. This is followed by a software-based trigger that reduces the accepted event rate to 1 kHz on average depending on the data-taking conditions.

3 Dataset and simulated events

The results were obtained using 139 fb^{-1} of pp LHC collision data collected between 2015 and 2018 by the ATLAS detector, with a centre-of-mass energy of 13 TeV and a 25 ns proton bunch crossing interval. In 2015–2016 the average number of interactions per bunch crossing (pile-up) was $\langle \mu \rangle = 20$, increasing to $\langle \mu \rangle = 38$ in 2017 and to $\langle \mu \rangle = 37$ in 2018. The uncertainty in the combined 2015–2018 integrated luminosity is 1.7% [31], obtained using the LUCID-2 detector [32] for the primary luminosity measurements.

Monte Carlo (MC) simulated datasets are used to model the SM backgrounds and evaluate signal selection efficiency and yields. All simulated samples were produced using the ATLAS simulation infrastructure [33] and GEANT 4 [34], or a faster simulation based on a parameterisation of the calorimeter response and GEANT 4 for the other detector systems. All simulated events were generated with a varying number of inelastic pp interactions overlaid on the hard-scattering event to model the multiple proton–proton interactions in the same and nearby bunch crossings. The pile-up events are generated with PYTHIA 8.186 [35] using the NNPDF2.3LO set of PDFs [36] and the A3 tune [37]. The simulated events were reconstructed with the same algorithms as those used for data.

The backgrounds considered in this analysis are: $t\bar{t}$ pair production; single-top production (s -channel, t -channel, and associated Wt production); W/Z +jets production; $t\bar{t}$ production with an electroweak boson ($t\bar{t}V$); Higgs boson pro-

Table 1 Overview of MC generators used for different simulated event samples

Process	Generator	Parton shower and hadronisation	Tune	PDF	Cross-section
$t\bar{t}$	POWHEG- BOX v2 [55–58]	PYTHIA 8.230 [35]	A14 [49]	NNPDF2.3LO [36]	NNLO+NNLL [59]
Single top	POWHEG- BOX v2 [60–62]	PYTHIA 8.230	A14	NNPDF2.3LO	NLO+NNLL [63]
W/Z +jets	SHERPA 2.2.1 [64]	SHERPA 2.2.1	SHERPA standard	NNPDF3.0NNLO	NNLO [65]
Diboson	SHERPA 2.2.1 & 2.2.2	SHERPA 2.2.1 & 2.2.2	SHERPA standard	NNPDF3.0NNLO	NLO
Triboson	SHERPA 2.2.1 & 2.2.2	SHERPA 2.2.1 & 2.2.2	SHERPA standard	NNPDF3.0NNLO	NLO
$t\bar{t} + V$	MADGRAPH5_aMC@NLO v2.3.3	PYTHIA 8.210	A14	NNPDF2.3LO	NLO [66]
tth	POWHEG- BOX v2	PYTHIA 8.230	AZNLO [67]	CTEQ6L1 [68]	NLO [69]
Vh	POWHEG- BOX v2	PYTHIA 8.212	A14	NNPDF2.3LO	NLO [69]

duction (tth , Vh); and diboson (WW , WZ , ZZ) and triboson (VVV where $V = W, Z$) production. Background samples were simulated using different MC event generators depending on the process. All background processes were normalised to the best available theoretical calculation of their respective cross-sections. The SHERPA samples used for W +jets modelling include up to two partons at NLO and four partons at LO using Comix [38] and OpenLoops [39, 40] and merged with the SHERPA parton shower [41] according to the ME+PS@NLO prescription [42–45] using the set of tuned parameters developed by the SHERPA authors. The event generators, the parton shower and hadronisation routines, and the underlying-event parameter tunes and parton distribution function (PDF) sets used in simulating the SM background processes, along with the accuracy of the theoretical cross-sections, are all summarised in Table 1.

For all samples showered with PYTHIA, the EvtGen v1.2.0 [46] program was used to simulate the properties of the bottom- and charm-hadron decays. Several samples produced without detector simulation were employed to estimate systematic uncertainties associated with the specific configuration of the MC generators used for the nominal SM background samples. They include variations of the renormalisation and factorisation scales, the CKKW-L [47] matching scale, as well as different PDF sets and fragmentation/hadronisation models. Details of the MC modelling uncertainties are discussed in Sect. 7.

The SUSY signal samples were generated using MADGRAPH5_aMC@NLO v2.6.2 [48] and PYTHIA 8.230 with the A14 [49] set of tuned parameters for the modelling of the parton showering (PS), hadronisation and underlying event. The matrix element (ME) calculation is performed at tree level and include the emission of up to two additional partons. The ME–PS matching is done using the CKKW-L prescription, with a matching scale set to one quarter of the chargino and next-to-lightest neutralino mass. The NNPDF2.3LO [36] PDF set was used.

Signal cross-sections are calculated at next-to-leading-order (NLO) accuracy in the strong coupling constant,

adding the resummation of soft gluon emission at next-to-leading-logarithm accuracy (NLO+NLL) [50–53]. The nominal cross-section and its uncertainty are taken as the midpoint and half-width of an envelope of cross-section predictions using different PDF sets and factorisation and renormalisation scales, as described in Ref. [54]. The simplified model has two parameters, the first being the mass of the $\tilde{\chi}_1^\pm$ and $\tilde{\chi}_2^0$ (which are assumed to be equal), and the second being the mass of the $\tilde{\chi}_1^0$. The signal cross-sections decrease as the $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ mass increases, ranging from 769 fb for a 250 GeV $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ mass to 1.3 fb for a 1000 GeV $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ mass.

4 Event reconstruction

Events are required to have at least one reconstructed interaction vertex with a minimum of two associated tracks each having $p_T > 500$ MeV. In events with multiple vertices, the one with the highest sum of squared transverse momenta of associated tracks is chosen as the primary vertex (PV) [70]. A set of baseline quality criteria are applied to reject events with non-collision backgrounds or detector noise [71].

Two identification levels are defined for leptons and jets: ‘baseline’ and ‘signal’. Baseline leptons and jets are selected with looser identification criteria, and are used in computing the missing transverse momentum as well as in resolving possible reconstruction ambiguities. Signal leptons and jets are a subset of the baseline objects with tighter quality requirements which are used to define the search regions. Isolation criteria, defined with a list of tracking-based and calorimeter-based variables, are used to select signal leptons by discriminating against semileptonic heavy-flavour decays and jets misidentified as leptons.

Electron candidates are reconstructed from energy deposits in the electromagnetic calorimeter that are matched to charged-particle tracks in the inner detector (ID) [72]. Baseline electrons are required to satisfy $p_T > 7$ GeV and $|\eta| < 2.47$. They are identified using the ‘loose’ operating point provided by a likelihood-based algorithm, described in

Ref. [72]. The number of hits in the innermost pixel layer is used to discriminate between electrons and converted photons. The longitudinal impact parameter z_0 relative to the PV is required to satisfy $|z_0 \sin \theta| < 0.5$ mm. The ‘tight’ likelihood operating point is applied for signal electron identification and the significance of the transverse impact parameter d_0 must satisfy $|d_0/\sigma(d_0)| < 5$. Signal electron candidates with $p_T < 200$ GeV are further refined using the *FCLoose* isolation working point, while those with larger p_T are required to pass the *FCHighPtCaloOnly* isolation working point, as described in Ref. [72].

Muon candidates are reconstructed from matching tracks in the ID and muon spectrometer, refined through a global fit which uses the hits from both subdetectors [73]. Baseline muons must have $p_T > 6$ GeV and $|\eta| < 2.7$, and satisfy the ‘medium’ identification criteria. Similarly to electrons, the longitudinal impact parameter z_0 relative to the PV is required to satisfy $|z_0 \sin \theta| < 0.5$ mm. Signal muon candidates are further defined with tighter pseudorapidity and impact parameter requirements, $|\eta| < 2.5$ and $|d_0/\sigma(d_0)| < 3$. The *FCLoose* isolation working point is also required for signal muons [73].

Jets are reconstructed from three-dimensional topological energy clusters in the calorimeters using the anti- k_t algorithm [74] with a radius parameter $R = 0.4$ [75]. Baseline jets are selected in the region $|\eta| < 4.5$ and have $p_T > 20$ GeV. To suppress jets from pile-up interactions, the jets with $|\eta| < 2.8$ and $p_T < 120$ GeV are required to satisfy the ‘medium’ working point of the jet vertex tagger (JVT), a tagging algorithm that identifies jets originating from the PV using track information [76,77]. The selection of signal jets is further refined by requiring them to be in the region $|\eta| < 2.8$ and have $p_T > 30$ GeV.

Jets containing b -hadrons are identified as ‘ b -tagged’ using the MV2c10 algorithm, a multivariate discriminant based on the track impact parameters and displaced secondary vertices [78]. These b -tagged jets are reconstructed in the region $|\eta| < 2.5$ and have $p_T > 30$ GeV. The b -tagging working point provides an efficiency of 77% for jets containing b -hadrons in simulated $t\bar{t}$ events, with rejection rates of 110 and 4.9 for light-flavour jets and jets containing c -hadrons, respectively [79].

To resolve the reconstruction ambiguities between electrons, muons and jets, an overlap removal procedure is applied to baseline objects. First, any electron sharing the same ID track with a muon is rejected. If it shares the same ID track with another electron, the one with lower p_T is discarded. Next, jets are rejected if they lie within $\Delta R = 0.2$ of a muon or if the muon is matched to the jet through ghost association [80]. Subsequently, electrons within a cone of size $\Delta R = \min(0.4, 0.04 + 10 \text{ GeV}/p_T)$ around a jet are removed. Last, muons within a cone, defined in the same way as for electrons, around any remaining jet are removed.

The missing transverse momentum $\mathbf{p}_T^{\text{miss}}$, with magnitude E_T^{miss} is calculated as the negative vectorial sum of the transverse momentum of all baseline reconstructed objects (electrons, muons, jets and photons [81]) and the soft term. The soft term includes all tracks associated with the PV but not matched to any reconstructed physics object. Tracks not associated with the PV are not considered in the E_T^{miss} calculation, improving the E_T^{miss} resolution by suppressing the effect of pile-up [82,83].

Corrections are applied to simulated events in order to account for the trigger, particle identification, and reconstruction efficiency differences between data and simulation.

5 Event selection

Events are recorded with the lowest-threshold E_T^{miss} trigger available, which is fully efficient for selecting events when the offline requirement of $E_T^{\text{miss}} > 240$ GeV is applied. To target the signal events, which have a leptonically decaying W boson and a Higgs boson decaying into a $b\bar{b}$ pair, events are required to have exactly one signal electron or muon (but not both) and either two or three signal jets, two of which must be b -tagged. The signal regions (SR) are defined using variables which suppress background contributions and increase the sensitivity for signal. These variables are based on the kinematic properties of the b -jets, the lepton and the missing transverse momentum, and are defined as follows:

- The invariant mass of the two b -jets, $m_{b\bar{b}}$, is required to be in the range $100 < m_{b\bar{b}} < 140$ GeV, in order to preferentially select b -jets from the Higgs boson decays.
- The invariant mass of the lepton and the leading b -jet is denoted by $m(\ell, b_1)$. For $t\bar{t}$ or single-top (particularly the Wt -channel) backgrounds, if the lepton and the leading b -jet originate from the same top-quark, the $m(\ell, b_1)$ distribution has an endpoint at $\sqrt{m^2(t) - m^2(W)}$. For signal events, the lepton and b -jet are produced from the $\tilde{\chi}_1^\pm$ and $\tilde{\chi}_2^0$ decay chains, respectively. The distribution of the invariant mass depends on the mass of the SUSY particles. For signal events with high-mass $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$, this observable provides good discrimination against background events.
- The transverse mass, m_T , is defined from the lepton transverse momentum \mathbf{p}_T^ℓ and $\mathbf{p}_T^{\text{miss}}$ as

$$m_T = \sqrt{2p_T^\ell E_T^{\text{miss}}(1 - \cos[\Delta\phi(\mathbf{p}_T^\ell, \mathbf{p}_T^{\text{miss}})])},$$

where $\Delta\phi(\mathbf{p}_T^\ell, \mathbf{p}_T^{\text{miss}})$ is the azimuthal angle between \mathbf{p}_T^ℓ and $\mathbf{p}_T^{\text{miss}}$. For W +jets and semileptonic $t\bar{t}$ events, in which one on-shell W boson decays leptonically, the observable has an upper endpoint at the W boson mass.

Table 2 Overview of the selection criteria for the signal regions. Each of the three ‘excl.’ SRs is binned in three m_{CT} regions for a total of nine ‘excl.’ bins

	SR-LM	SR-MM	SR-HM
N_{lepton}		= 1	
p_T^ℓ [GeV]		> 7(6) for $e(\mu)$	
N_{jet}		= 2 or 3	
$N_{b\text{-jet}}$		= 2	
E_T^{miss} [GeV]		> 240	
$m_{b\bar{b}}$ [GeV]		∈ [100, 140]	
$m(\ell, b_1)$ [GeV]	–	–	> 120
m_T [GeV] (excl.)	∈ [100, 160]	∈ [160, 240]	> 240
m_{CT} [GeV] (excl.)		{ ∈ [180, 230], ∈ [230, 280], > 280 }	
m_T [GeV] (disc.)	> 100	> 160	> 240
m_{CT} [GeV] (disc.)		> 180	

The m_T distribution for signal events extends significantly above $m(W)$.

- The contranverse mass [84,85] of two b -jets, m_{CT} , is defined as:

$$m_{CT} = \sqrt{2p_T^{b_1} p_T^{b_2} (1 + \cos \Delta\phi_{bb})},$$

where $p_T^{b_1}$ and $p_T^{b_2}$ are the transverse momenta of the two leading b -jets and $\Delta\phi_{bb}$ is the azimuthal angle between them. For the $t\bar{t}$ background, the observable has an upper endpoint at $(m^2(t) - m^2(W))/m(t)$. A requirement that m_{CT} be larger than 180 GeV efficiently suppresses the $t\bar{t}$ background.

An overview of the signal region definitions is provided in Table 2. Three separate classes of signal regions are defined, progressively targeting increasing mass differences between the $\tilde{\chi}_1^\pm$ (and its mass-degenerate $\tilde{\chi}_2^0$ wino partner) and the $\tilde{\chi}_1^0$. These regions are labelled SR-LM, SR-MM and SR-HM to indicate low (LM), medium (MM) and high (HM) mass differences respectively. Requirements on m_T make the three regions mutually exclusive. Of the three signal regions, SR-LM selects the smallest values of m_T . It targets signal models with a mass-splitting between the $\tilde{\chi}_2^0$ (and hence the $\tilde{\chi}_1^\pm$) and the $\tilde{\chi}_1^0$ that is similar to the Higgs boson mass. The other two signal regions select progressively larger mass differences by requiring larger values of m_T . The signal region with the highest requirement on m_T , SR-HM, also requires $m(\ell, b_1) > 120$ GeV in order to further suppress $t\bar{t}$ and single-top background events. The three signal regions otherwise share a common set of selections on E_T^{miss} , $m_{b\bar{b}}$ and m_{CT} .

When setting model-dependent exclusion limits (‘excl.’), each of the three SRs is binned in three m_{CT} regions, thus providing nine bins in total for a simultaneous two-dimensional fit in m_{CT} and m_T across the three SRs. This multi-bin

approach enhances the sensitivity to a range of SUSY scenarios with different properties. For model-independent limits and null-hypothesis tests (‘disc.’ for discovery), the various m_{CT} bins are merged for each of the three SRs. The requirement of $m(\ell, b_1) > 120$ GeV is only applied in SR-HM. Furthermore, the upper bound on m_T is removed for SR-LM and SR-MM. The fit strategy is detailed in Sect. 6. The systematic uncertainties, fit and results discussed in the following sections are based on the exclusion SRs, while the model-independent results are based on the discovery SRs.

6 Background estimation

The expected backgrounds in each signal region are determined in a profile likelihood fit, referred to as a ‘background-only fit’. In this fit, the normalisation of the backgrounds is adjusted to match the data in control regions with negligible signal contamination. The resulting normalisation factors are then used to correct the expected yields of the corresponding backgrounds in the various signal regions. The control regions – as detailed in the following – are designed to be enriched in the major background processes: $t\bar{t}$, single-top and W +jets processes. The control region for single-top has a similar composition in single-top processes as the signal regions, therefore a single scale factor is used. All CRs are designed to be non-overlapping with the signal regions and also mutually exclusive. A probability density function is defined for each of the control regions. The inputs are the observed event yield and the predicted background yield from simulation with Poisson statistical uncertainties as well as with systematic uncertainties as nuisance parameters. The nuisance parameters are constrained by Gaussian distributions with widths corresponding to the sizes of the uncertainties. The uncertainties do not only vary the scales, but also account for bin-to-bin transitions. The systematic uncertainties are detailed in Sect. 7. The prod-

Table 3 Overview of the CR and VR definitions. All regions partially share the same selection as the SR for all variables except $m(\ell, b_1)$, which is not used in the CR and VR definitions

CR	TR-LM	TR-MM	TR-HM	WR	STR	
$m_{b\bar{b}}$ [GeV]		<100 or >140		$\in [50, 80]$	>195	
m_T [GeV]	$\in [100, 160]$	$\in [160, 240]$	>240	$\in [50, 100]$	>100	
m_{CT} [GeV]		<180		>180	>180	
VR	VR-onLM	VR-onMM	VR-onHM	VR-offLM	VR-offMM	VR-offHM
$m_{b\bar{b}}$ [GeV]		$\in [100, 140]$		$\in [50, 80] \cup [160, 195]$	$\in [50, 80] \cup [160, 195]$	$\in [50, 75] \cup [165, 195]$
m_T [GeV]	$\in [100, 160]$	$\in [160, 240]$	>240	$\in [100, 160]$	$\in [160, 240]$	>240
m_{CT} [GeV]		<180			>180	

uct of all the probability density functions forms the likelihood. Normalisation and nuisance parameters are correlated in all regions participating in the fit. The likelihood is maximised by adjusting the normalisation and nuisance parameters. The extrapolation of the adjusted normalisation and nuisance parameters to the signal regions is checked in validation regions (VR), as defined below, which kinematically resemble the signal regions but are expected to have less signal. The VRs do not overlap with the CRs or SRs.

Subdominant background processes, such as Z +jets, diboson and multiboson, $t\bar{t}+V$, $t\bar{t}+h$ and Vh , which have no dedicated control regions, are normalised to the cross-sections indicated in Table 1. In the same way as for the dominant backgrounds, their expected yields in the SRs are subject to statistical and systematic uncertainties. Backgrounds with fake leptons such as jets misreconstructed as a lepton, and events with leptons originating from a jet produced by heavy-flavour quarks or from photon conversions are estimated using a matrix method as described in Ref. [86], and found to be negligible in all regions.

The $t\bar{t}$ background estimation relies on a set of three CRs (labelled TR-LM, TR-MM, TR-HM), each with an m_T selection the same as in the SRs. In order to obtain samples enriched in $t\bar{t}$ events, the requirement on $m(\ell, b_1)$ is removed and the selection criteria for m_{CT} and $m_{b\bar{b}}$ are inverted relative to the SRs. These three control regions are fit simultaneously to obtain a single normalisation factor. The W +jets contributions in the SRs are constrained by a single CR (labelled WR) defined similarly to the SRs but with less stringent lower bounds on m_T and an off-peak region for $m_{b\bar{b}}$. The fraction of events with heavy flavor hadrons in simulated W +jet events was found similar between the WR and the SRs. Events in the single-top CR (labelled STR) must satisfy the SR requirements except that this CR requires $m_T > 100$ GeV and $m_{b\bar{b}} > 195$ GeV. The $t\bar{t}$ purity varies from 79% in TR-LM to 86% in TR-MM.

The purity of the single-top (W +jets) is 52% (53%) in STR (WR).

Two sets of VRs are defined for each SR, including the off-peak ($m_{b\bar{b}} < 100$ or > 140 GeV) and the on-peak $m_{b\bar{b}}$ regions, with the same m_T as in the SR. The on-peak VRs validate the extrapolation from the CRs to the SRs in $m_{b\bar{b}}$, and the off-peak VRs validate the extrapolation in m_{CT} . The validation regions share the same m_T binning as the signal regions, denoted as LM, MM and HM. The background modelings are validated in low, medium and high m_T regions separately. A summary of all CR and VR selection criteria is reported in Table 3.

7 Systematic uncertainties

Systematic uncertainties are evaluated for all simulated signal and background events. For the dominant backgrounds with dedicated control regions, the systematic uncertainties impact the extrapolation from the control regions to the corresponding signal regions. For all other backgrounds estimated from simulation, the uncertainties affect the overall cross-section normalisation and the acceptance of the analysis selection. Uncertainties arising from theoretical modelling and detector effects are estimated and discussed below. A breakdown of the dominant systematic uncertainties in background estimates in the various exclusion signal regions is summarised in Table 4. The uncertainties in the scale factor fits to the control regions are listed as ‘Normalisation of dominant backgrounds’.

Several uncertainties in the theoretical modelling of the single-top and $t\bar{t}$ backgrounds are considered. Uncertainties due to the choice of hard-scatter generation program are estimated by comparing POWHEG-BOX generated events, showered using PYTHIA 8, with events generated by aMC@NLO and showered with PYTHIA 8, while those due to the choice of parton shower model are evaluated by comparing POWHEG-BOX generated samples showered using PYTHIA 8 with

Table 4 Breakdown of the dominant systematic uncertainties in background estimates in the various exclusion signal regions. The individual uncertainties can be correlated, and do not necessarily add up in quadrature to the total background uncertainty. The percentages show the size of the uncertainty relative to the total expected background

Signal region	SR-LM	SR-MM	SR-HM
Total background expectation	27	8.6	8.1
Total uncertainty	±4 [15%]	±2.2 [25%]	±2.7 [34%]
Theoretical systematic uncertainties			
$t\bar{t}$	±2.6 [10%]	±0.6 [7%]	±0.33 [4%]
Single top	±0.8 [2.7%]	±1.1 [12%]	±1.9 [23%]
W+jets	±0.23 [0.9%]	±0.07 [0.8%]	±0.19 [2.3%]
Other backgrounds	±0.13 [0.5%]	±0.15 [1.7%]	±0.08 [1.0%]
MC statistical uncertainties			
MC statistics	±1.7 [6%]	±1.1 [13%]	±1.2 [14%]
Uncertainties in the background normalisation			
Normalisation of dominant backgrounds	±1.3 [5%]	±1.6 [18%]	±1.3 [16%]
Experimental systematic uncertainties			
$E_T^{\text{miss}}/JVT/\text{pile-up}/\text{trigger}$	±1.8 [7%]	±0.4 [4%]	±0.4 [5%]
Jet energy resolution	±1.6 [6%]	±0.5 [6%]	±0.4 [5%]
b-tagging	±1.1 [4%]	±0.29 [3.4%]	±0.13 [1.5%]
Jet energy scale	±0.9 [3.2%]	±0.9 [10%]	±0.29 [4%]
Lepton uncertainties	±0.32 [1.2%]	±0.09 [1.0%]	±0.19 [2.3%]

Table 5 Background fit results for the exclusion SR regions. The errors shown are the statistical plus systematic uncertainties. Uncertainties in the fitted yields are symmetric by construction, except where the negative error is truncated at an event yield of zero

	All m_{CT} bins	Low m_{CT}	Medium m_{CT}	High m_{CT}
SR-LM				
Observed	34	16	11	7
Expected	27 ± 4	8.8 ± 2.8	11.3 ± 3.1	7.3 ± 1.5
$t\bar{t}$	16.2 ± 3.4	4.4 ± 2.2	7.3 ± 2.5	4.6 ± 1.2
Single top	2.7 ± 1.8	1.3 ± 1.1	0.9 ^{+1.0} _{-0.9}	0.6 ± 0.6
W+jets	5.5 ± 2.0	2.0 ± 0.9	2.4 ± 1.3	1.1 ± 0.5
Di-/Multiboson	0.67 ± 0.19	0.39 ± 0.13	0.09 ^{+0.11} _{-0.09}	0.18 ± 0.04
Others	2.23 ± 0.29	0.81 ± 0.25	0.64 ± 0.15	0.77 ± 0.12
SR-MM				
Observed	13	4	7	2
Expected	8.6 ± 2.2	4.6 ± 1.7	2.6 ± 1.3	1.4 ± 0.6
$t\bar{t}$	2.7 ± 1.4	1.6 ± 0.9	0.8 ± 0.7	0.30 ± 0.24
Single top	2.7 ± 1.9	1.6 ± 1.5	1.0 ^{+1.1} _{-1.0}	0.15 ^{+0.19} _{-0.15}
W+jets	1.5 ± 0.7	0.6 ± 0.4	0.3 ^{+0.4} _{-0.3}	0.57 ± 0.26
Di-/Multiboson	0.29 ± 0.08	0.09 ± 0.04	0.065 ± 0.028	0.14 ± 0.06
Others	1.33 ± 0.27	0.69 ± 0.20	0.40 ± 0.13	0.24 ± 0.09
SR-HM				
Observed	14	6	5	3
Expected	8.1 ± 2.7	4.1 ± 1.9	2.9 ± 1.3	1.1 ± 0.5
$t\bar{t}$	1.4 ± 0.5	0.8 ± 0.4	0.36 ± 0.25	0.22 ± 0.15
Single top	2.0 ^{+2.4} _{-2.0}	0.9 ^{+1.5} _{-0.9}	0.9 ± 0.9	0.16 ^{+0.26} _{-0.16}
W+jets	3.7 ± 1.0	1.9 ± 0.8	1.4 ± 0.8	0.45 ± 0.19
Di-/Multiboson	0.21 ± 0.06	0.057 ± 0.025	0.075 ± 0.027	0.08 ± 0.04
Others	0.74 ± 0.16	0.34 ± 0.09	0.19 ± 0.08	0.21 ± 0.08

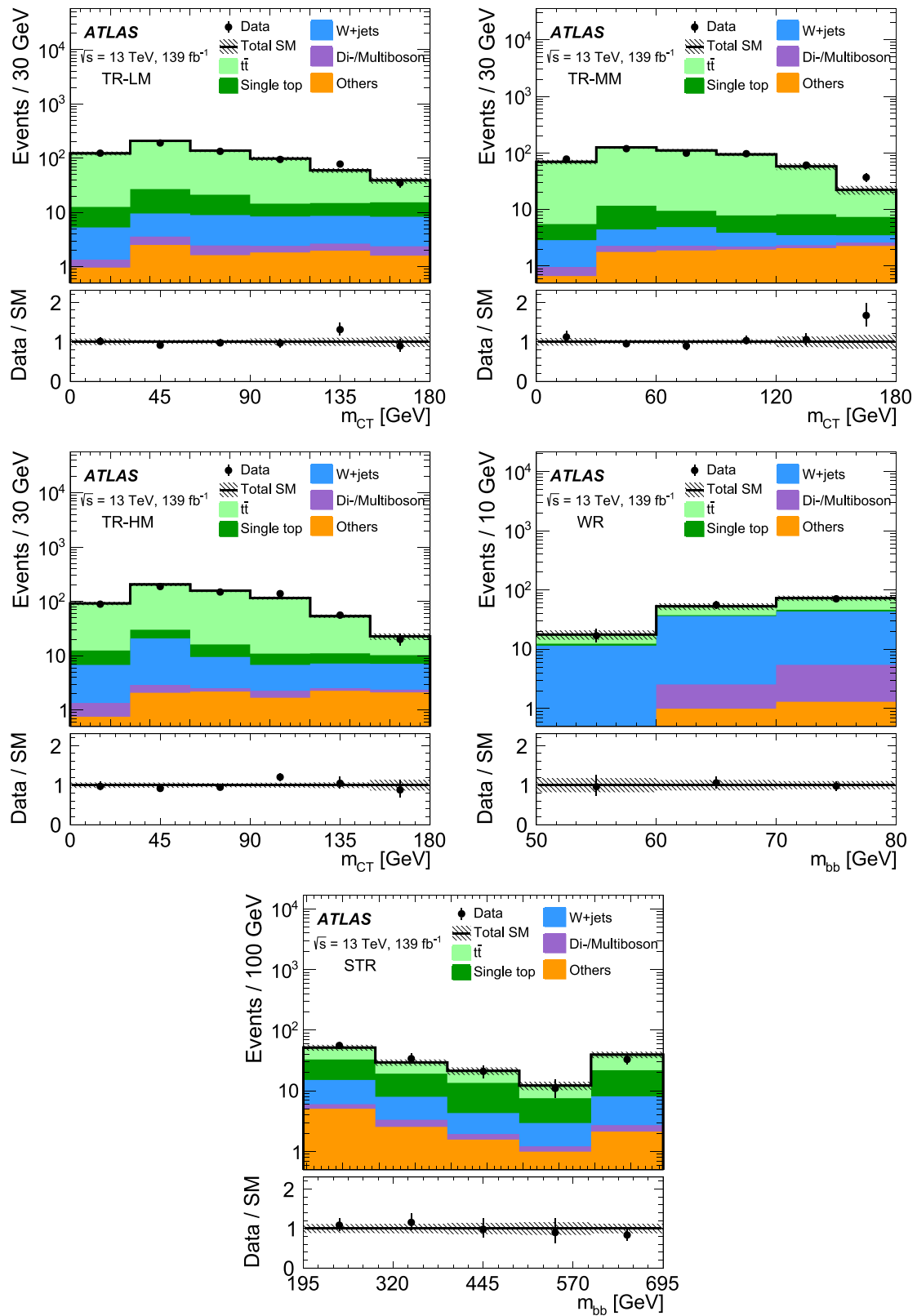


Fig. 2 The post-fit m_{CT} distributions in TR-LM, TR-MM, and TR-HM are shown as well as the post-fit $m_{b\bar{b}}$ distributions in WR and STR. The uncertainty bands plotted include all statistical and systematic uncertainties. The overflow events, where present, are included in the last bin

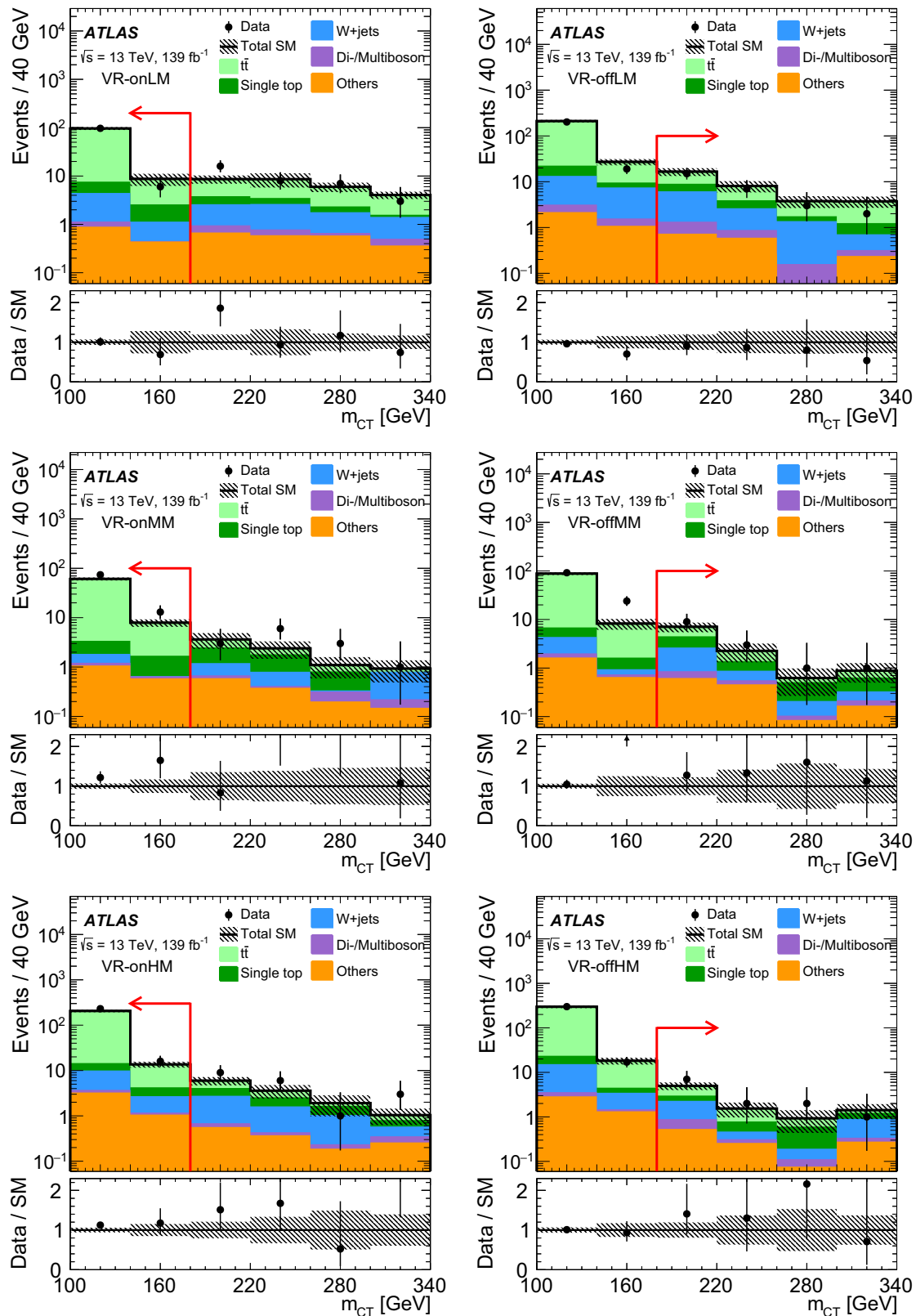


Fig. 3 The post-fit m_{CT} distributions are shown in each of the validation regions (VR-onLM, VR-onMM, VR-onHM, VR-offLM, VR-offMM, and VR-offHM) after all the selection requirements are applied other than the m_{CT} cut. The uncertainty bands plotted include all sta-

tistical and systematic uncertainties. The overflow (underflow) events, where present, are included in the last (first) bin. The line with an arrow indicates the m_{CT} cut used in VR selection

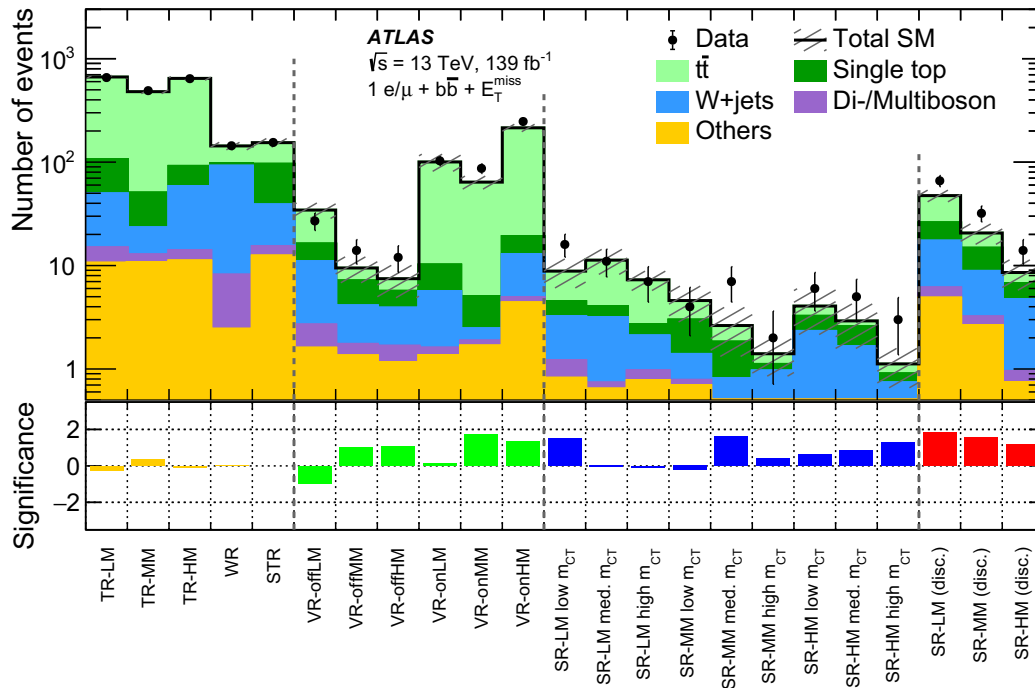


Fig. 4 Comparison of the observed and expected event yields in control, validation, exclusion, and discovery signal regions. Uncertainties in the background estimates include both the statistical (in the simulated

event yields) and systematic uncertainties. The bottom panel shows the significance [91] of the differences between the observed and expected yields. Not all regions shown here are statistically independent

POWHEG-BOX samples showered using HERWIG 7 [87]. The uncertainties from the modelling of initial- and final-state radiation are assessed by varying the renormalisation and factorisation scales up and down by a factor of two, with the radiation setting varied as well [88]. For single-top Wt production, the impact of interference between single-resonant and double-resonant top-quark production is estimated by comparing the nominal sample generated using the diagram removal method with a sample generated using the diagram subtraction method. For the different signal regions, the dominant uncertainty sources are the $t\bar{t}$ parton shower in SR-LM (10%), and the single-top generator uncertainties in SR-MM (10%) and SR-HM (21%).

Theory uncertainties affecting the generator predictions for W/Z +jets, diboson, triboson and $t\bar{t} + W/Z$ samples are estimated by taking the envelope of the seven-point variations of the renormalisation and factorisation scales. For W/Z +jets, the uncertainties from the PDF variations, as well as from the variations of matching and resummation scales are also considered. Additionally, an overall 6% (20%) systematic uncertainty in the inclusive cross-section is assigned for the diboson (triboson) sample [89]. Similar cross-section uncertainties are also assigned for other small background contributions.

Theory uncertainties in the expected yields for SUSY signals are estimated by varying by a factor of two the parameters corresponding to the factorisation, renormalisation, and CKKW-L matching scales, as well as the PYTHIA 8 shower tune parameters. The overall uncertainties range from about 10% in the region with a large splitting between the $\tilde{\chi}_2^0/\tilde{\chi}_1^\pm$ and $\tilde{\chi}_1^0$ masses to about 25% in the mass spectra with small mass splitting.

The dominant detector systematic effects are the uncertainties associated with the jet energy scale (JES) and jet energy resolution (JER), the E_T^{miss} modelling, and pile-up. The jet uncertainties are measured as a function of the p_T and η of the jet, the pile-up conditions and the jet flavour composition. They are determined using a combination of data and simulation, through measurements of the jet p_T balance in dijet, Z +jets and γ +jets events [90]. The systematic uncertainties in the E_T^{miss} modelling are derived by propagating the uncertainties in the energy and momentum scale of each of the objects entering the calculation, and the uncertainties in the soft term's resolution and scale [83]. A pile-up reweighting procedure is applied to simulation to match the distribution of the number of reconstructed vertices observed in data. The corresponding uncertainty is derived by a reweighting in which $\langle\mu\rangle$ is varied by $\pm 4\%$. The experimental uncertainties have a less significant impact than the theoretical ones in all

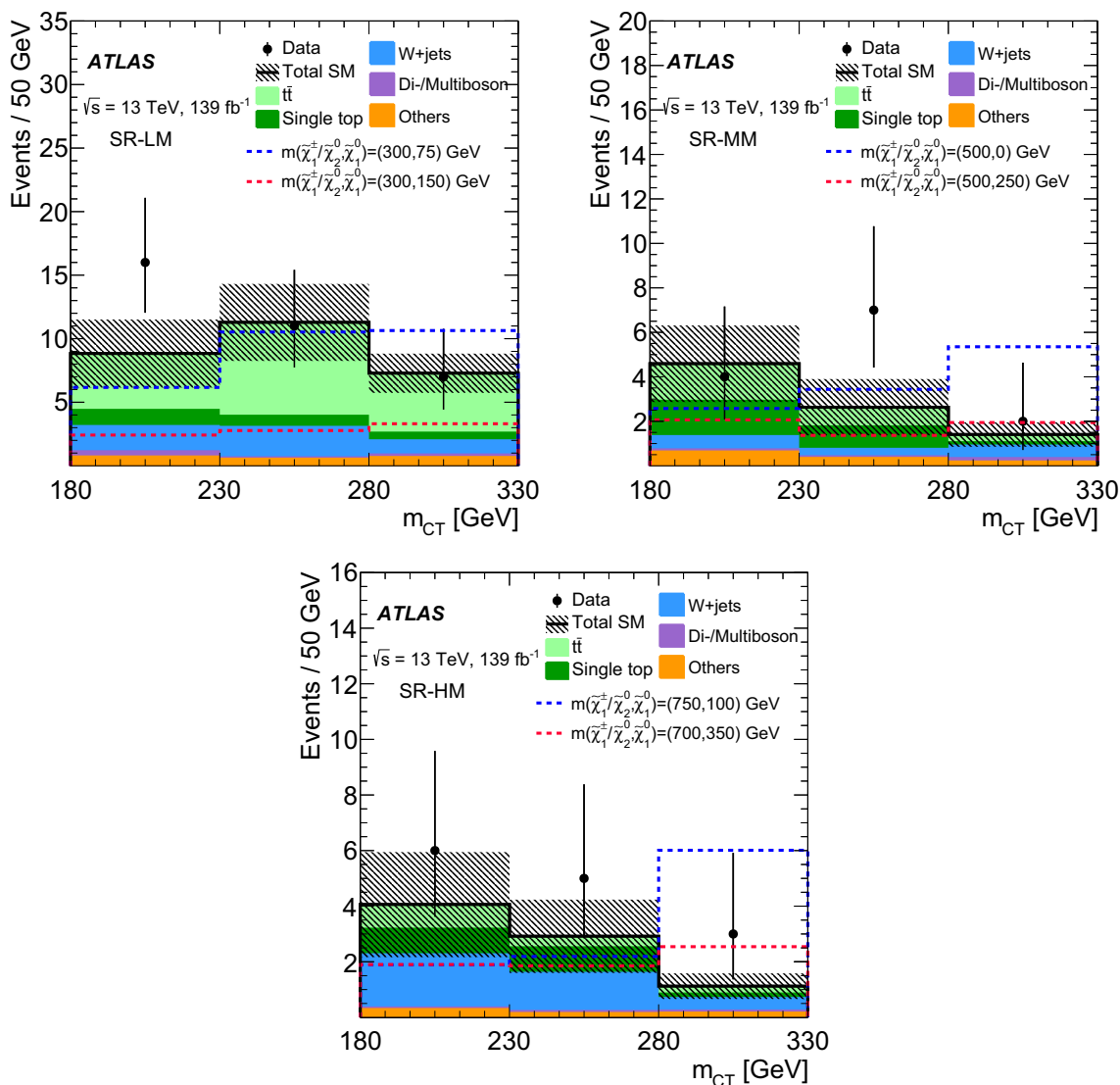


Fig. 5 The post-fit m_{CT} distributions in the exclusion signal regions (SR-LM, SR-MM, and SR-HM). The uncertainty bands plotted include all statistical and systematic uncertainties. The dashed lines represent the benchmark signal samples. The overflow events, where present, are included in the last bin

signal regions: the largest experimental source contributes 5–10% depending on the SR. The MC statistical uncertainties contribute 5–18% depending on the SR.

8 Results

The observed event yield in each of the exclusion signal regions is summarised in Table 5 along with the corresponding Standard Model predictions obtained from the background-only fit. The background normalisation factors are $1.02^{+0.07}_{-0.09}$ for $t\bar{t}$, $0.6^{+0.5}_{-0.25}$ for single top, and $1.22^{+0.26}_{-0.24}$ for W+jets.

In Fig. 2 the post-fit m_{CT} distributions in the $t\bar{t}$ control regions TR-LM, TR-MM, and TR-HM are compared with the data. For the W boson and single-top control regions the $m_{b\bar{b}}$ distribution is shown. Figure 3 shows the post-fit m_{CT} distributions after all of the validation region selection requirements are applied except the m_{CT} cut. The data and the background expectation in all validation regions agree well within two standard deviations. Therefore no further systematic uncertainty is applied on the background estimation in the signal regions.

The compatibility of the observed and expected event yields in control, validation, exclusion, and discovery signal regions is illustrated in Fig. 4. No significant excess over

Table 6 Left to right: 95% CL upper limits on the visible cross-section ($\langle\epsilon\sigma\rangle_{\text{obs}}^{95}$) and on the number of signal events (S_{obs}^{95}). The third column (S_{exp}^{95}) shows the expected 95% CL upper limit (and its $\pm 1\sigma$ excursions) on the number of signal events if no BSM signal is present. The last

three columns indicate the CL_B value, i.e. the confidence level observed for the background-only hypothesis, the discovery p -value (p_0) and the significance Z [91]

Signal Region	$\langle\epsilon\sigma\rangle_{\text{obs}}^{95}$ [fb]	S_{obs}^{95}	S_{exp}^{95}	CL_B	p_0	Z
SR-LM(disc.)	0.26	36.8	$20.0^{+8.0}_{-5.4}$	0.97	0.03	1.88
SR-MM(disc.)	0.18	24.8	$15.3^{+6.2}_{-4.6}$	0.94	0.06	1.54
SR-HM(disc.)	0.11	14.7	$9.7^{+3.3}_{-2.7}$	0.89	0.10	1.30

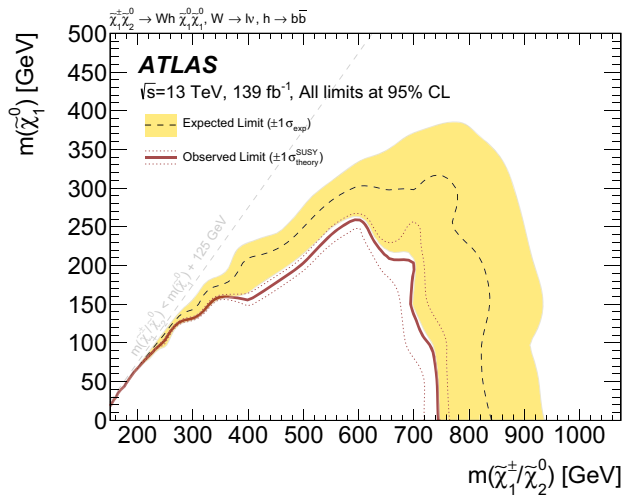


Fig. 6 Model-dependent exclusion contour at 95% CL on the production of a chargino and a next-to-lightest neutralino. The observed limit is given by the solid line with the signal cross-section uncertainties shown by the dotted lines as indicated in the text. Expected limits are given by the dashed line with uncertainties shown by the shaded band

the SM prediction is observed in data. Figure 5 shows the post-fit m_{CT} distributions in SR-LM, SR-MM, and SR-HM. The uncertainty bands include all statistical and systematic uncertainties. The dashed lines represent the benchmark signal points.

Model-dependent exclusion limits at 95% confidence level (CL) are placed on the signal model. These limits are shown as a function of the masses of the supersymmetric particles in Fig. 6. They are determined interpolating between the simulated mass points, but no smoothing procedure is applied so that local fluctuations in the limit can be present. A likelihood similar to the one used in the background-only fit, but with additional terms for the SRs, is used for the calculation. The exclusion SRs thus participate in the fit and are used to constrain normalisation and nuisance parameters. A signal is allowed in this likelihood in both the CRs and SRs. The VRs are not used in the fit. The CL_s method [92] is used to derive the confidence level of the exclusion for a particular signal model; signal models with a CL_s value below 0.05 are excluded at 95% CL. The uncertainties in the observed

limit are calculated by varying the cross-section for the signal up and down by its uncertainty. Due to a modest excess observed in some bins of the exclusion signal regions, the observed limit is weaker than the expected limit and extends up to about 740 GeV in $m(\tilde{\chi}_1^\pm/\tilde{\chi}_2^0)$ for massless $\tilde{\chi}_1^0$ and up to $m(\tilde{\chi}_1^\pm/\tilde{\chi}_2^0) = 600$ GeV for $m(\tilde{\chi}_1^0) = 250$ GeV. Benefiting from the increased integrated luminosity and the improved two-dimensional fit procedure, the current observed limit exceeds the previous ATLAS limit by about 200 GeV in $m(\tilde{\chi}_1^\pm/\tilde{\chi}_2^0)$ for a massless $\tilde{\chi}_1^0$.

Table 6 summarises the observed (S_{obs}) and expected (S_{exp}) 95% CL upper limits on the number of signal events and on the observed visible cross-section, $\langle\epsilon\sigma\rangle_{\text{obs}}^{95}$, for each of the three cumulative discovery signal regions. These cumulative signal regions are those defined to test for the presence of any beyond-the-Standard-Model (BSM) physics processes, where in every case the upper bound on m_T is also removed. Upper limits on contributions from new physics processes are estimated using the so-called ‘model-independent fit’, where a generic BSM process is assumed to contribute only to the SR and not to the CRs, thus giving a conservative background estimate in the SR. When normalised to the integrated luminosity of the data sample, the results can be interpreted as corresponding to observed upper limits $\langle\epsilon\sigma\rangle_{\text{obs}}^{95}$, defined as the product of the production cross-section, the acceptance, and the selection efficiency of a BSM signal. The p_0 values, which represent the probability of the SM background alone to fluctuate to the observed number of events or higher, are also provided. All numbers are calculated from pseudo-experiments.

9 Conclusion

The results of a search for electroweakino pair production $pp \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_2^0$ in which the chargino ($\tilde{\chi}_1^\pm$) decays into a W boson and the lightest neutralino ($\tilde{\chi}_1^0$), while the heavier neutralino ($\tilde{\chi}_2^0$) decays into the Standard Model 125 GeV Higgs boson and a second $\tilde{\chi}_1^0$ are presented. The analysis is performed using pp collisions provided by the LHC at a centre-of-mass energy of 13 TeV. Data collected with the ATLAS detector between 2015 and 2018 are used, corresponding to

an integrated luminosity of 139 fb^{-1} . No significant deviation from the expected Standard Model background is observed. Limits are set on the direct production of the electroweakino in simplified models. Masses of $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ up to 740 GeV are excluded at 95% confidence level for a massless $\tilde{\chi}_1^0$. The current search improves on the previous ATLAS limit by about 200 GeV in $m(\tilde{\chi}_1^\pm/\tilde{\chi}_2^0)$ for a massless $\tilde{\chi}_1^0$.

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Data Availability Statement This manuscript has no associated data or the data will not be deposited. [Authors’ comment: All ATLAS scientific output is published in journals, and preliminary results are made available in Conference Notes. All are openly available, without restriction on use by external parties beyond copyright law and the standard conditions agreed by CERN. Data associated with journal publications are also made available: tables and data from plots (e.g. cross section values, likelihood profiles, selection efficiencies, cross section limits, ...) are stored in appropriate repositories such as HEPDATA (<http://hepdata.cedar.ac.uk/>). ATLAS also strives to make additional material related to the paper available that allows a reinterpretation of the data in the context of new theoretical models. For example, an extended encapsulation of the analysis is often provided for measurements in the framework of RIVET (<http://rivet.hepforge.org/>).]

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W. M. Barbe³⁸, E. L. Barberio¹⁰⁵, D. Barberis^{55a,55b}, M. Barbero¹⁰², G. Barbour⁹⁵, T. Barillari¹¹⁵, M. S. Barisits³⁶, J. Barkeloo¹³², T. Barklow¹⁵³, R. Barnea¹⁶⁰, S. L. Barnes^{60c}, B. M. Barnett¹⁴⁴, R. M. Barnett¹⁸, Z. Barnovska-Blenessy^{60a}, A. Baroncelli^{60a}, G. Barone²⁹, A. J. Barr¹³⁵, L. Barranco Navarro^{45a,45b}, F. Barreiro⁹⁹, J. Barreiro Guimarães da Costa^{15a}, S. Barsov¹³⁸, R. Bartoldus¹⁵³, G. Bartolini¹⁰², A. E. Barton⁹⁰, P. Bartos^{28a}, A. Basalae⁴⁶, A. Basan¹⁰⁰, A. Bassalat^{65.ap}, M. J. Basso¹⁶⁷, R. L. Bates⁵⁷, S. Batlamous^{35e}, J. R. Batley³², B. Batool¹⁵¹, M. Battaglia¹⁴⁶, M. Bauce^{73a,73b}, F. Bauer¹⁴⁵, K. T. Bauer¹⁷¹, H. S. Bawa^{31,1}, J. B. Beacham⁴⁹, T. Beau¹³⁶, P. H. Beauchemin¹⁷⁰, F. Becherer⁵², P. Bechtle²⁴, H. C. Beck⁵³, H. P. Beck^{20,r}, K. Becker⁵², M. Becker¹⁰⁰, C. Becot⁴⁶, A. Beddall^{12d}, A. J. Beddall^{12a}, V. A. Bednyakov⁸⁰, M. Bedognetti¹²⁰, C. P. Bee¹⁵⁵, T. A. Beermann¹⁸², M. Begalli^{81b}, M. Begel²⁹, A. Behera¹⁵⁵, J. K. Behr⁴⁶, F. 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Calvet¹⁵⁵, M. Calvetti^{72a,72b}, R. Camacho Toro¹³⁶, S. Camarda³⁶, D. Camarero Munoz⁹⁹, P. Camarri^{74a,74b}, D. Cameron¹³⁴, R. Caminal Armadans¹⁰³, C. Camincher³⁶, S. Campana³⁶, M. Campanelli⁹⁵, A. Camplani⁴⁰, A. Campoverde¹⁵¹, V. Canale^{70a,70b}, A. Canesse¹⁰⁴, M. Cano Bret^{60c}, J. Cantero¹³⁰, T. Cao¹⁶¹, Y. Cao¹⁷³, M. D. M. Capeans Garrido³⁶, M. Capua^{41a,41b}, R. Cardarelli^{74a}, F. Cardillo¹⁴⁹, G. Carducci^{41a,41b}, I. Carli¹⁴³, T. Carli³⁶, G. Carlino^{70a}, B. T. Carlson¹³⁹, L. Carminati^{69a,69b}, R. M. D. Carney^{45a,45b}, S. Caron¹¹⁹, E. Carquin^{147d}, S. Carrá⁴⁶, J. W. S. Carter¹⁶⁷, M. P. Casado^{14.e}, A. F. Casha¹⁶⁷, D. W. Casper¹⁷¹, R. Castelijin¹²⁰, F. L. Castillo¹⁷⁴, V. Castillo Gimenez¹⁷⁴, N. F. Castro^{140a,140e}, A. Catinaccio³⁶, J. R. Catmore¹³⁴, A. Cattai³⁶, J. Caudron²⁴, V. Cavaliere²⁹, E. Cavallaro¹⁴, M. Cavalli-Sforza¹⁴, V. Cavasinni^{72a,72b}, E. Celebi^{12b}, F. Ceradini^{75a,75b}, L. Cerda Alberich¹⁷⁴, K. Cerny¹³¹, A. S. Cerqueira^{81a}, A. Cerri¹⁵⁶, L. Cerrito^{74a,74b}, F. Cerutti¹⁸, A. Cervelli^{23a,23b}, S. A. Cetin^{12b}, Z. Chadi^{35a}, D. Chakraborty¹²¹, W. S. Chan¹²⁰, W. Y. Chan⁹¹, J. D. Chapman³², B. Chargeishvili^{159b}, D. G. Charlton²¹, T. P. Charman⁹³, C. C. Chau³⁴, S. Che¹²⁷, S. Chekanov⁶, S. V. Chekulaev^{168a}, G. A. Chelkov⁸⁰, M. A. Chelstowska³⁶, B. Chen⁷⁹, C. Chen^{60a}, C. H. Chen⁷⁹, H. Chen²⁹, J. Chen^{60a}, J. Chen³⁹, S. Chen¹³⁷, S. J. Chen^{15c}, X. Chen^{15b}, Y. Chen⁸³, Y. H. Chen⁴⁶, H. C. Cheng^{63a}, H. J. Cheng^{15a}, A. Cheplakov⁸⁰, E. Cheremushkina¹²³, R. Cherkaoui El Moursli^{35e}, E. Cheu⁷, K. Cheung⁶⁴, T. J. A. Chevaléras¹⁴⁵, L. Chevalier¹⁴⁵, V. Chiarella⁵¹, G. Chiarelli^{72a}, G. Chiodini^{68a}, A. S. Chisholm²¹, A. Chitan^{27b}, I. Chiu¹⁶³, Y. H. Chiu¹⁷⁶, M. V. Chizhov⁸⁰, K. Choi⁶⁶,

A. R. Chomont^{73a,73b}, S. Chouridou¹⁶², Y. S. Chow¹²⁰, M. C. Chu^{63a}, X. Chu^{15a,15d}, J. Chudoba¹⁴¹, A. J. Chuinard¹⁰⁴, J. J. Chwastowski⁸⁵, L. Chytka¹³¹, D. Cieri¹¹⁵, K. M. Ciesla⁸⁵, D. Cinca⁴⁷, V. Cindro⁹², I. A. Cioară^{27b}, A. Ciochio¹⁸, F. Ciroto^{70a,70b}, Z. H. Citron^{180j}, M. Citterio^{69a}, D. A. Ciubotaru^{27b}, B. M. Ciungu¹⁶⁷, A. Clark⁵⁴, M. R. Clark³⁹, P. J. Clark⁵⁰, C. Clement^{45a,45b}, Y. Coadou¹⁰², M. Cobl^{67a,67c}, A. Coccaro^{55b}, J. Cochran⁷⁹, H. Cohen¹⁶¹, A. E. C. Coimbra³⁶, L. Colasurdo¹¹⁹, B. Cole³⁹, A. P. Colijn¹²⁰, J. Collot⁵⁸, P. Conde Muño^{140a,140h}, E. Coniavitis⁵², S. H. Connell^{33c}, I. A. Connelly⁵⁷, S. Constantinescu^{27b}, F. Conventi^{70a,av}, A. M. Cooper-Sarkar¹³⁵, F. Cormier¹⁷⁵, K. J. R. Cormier¹⁶⁷, L. D. Corpe⁹⁵, M. Corradi^{73a,73b}, E. E. Corrigan⁹⁷, F. Corriveau^{104,ae}, A. Cortes-Gonzalez³⁶, M. J. Costa¹⁷⁴, F. Costanza⁵, D. Costanzo¹⁴⁹, G. Cowan⁹⁴, J. W. Cowley³², J. Crane¹⁰¹, K. Cranmer¹²⁵, S. J. Crawley⁵⁷, R. A. Creager¹³⁷, S. Crépe-Renaudin⁵⁸, F. Crescioli¹³⁶, M. Cristinziani²⁴, V. Croft¹²⁰, G. Crosetti^{41a,41b}, A. Cueto⁵, T. Cuhadar Donszelmann¹⁴⁹, A. R. Cukierman¹⁵³, W. R. Cunningham⁵⁷, S. Czekierda⁸⁵, P. Czodrowski³⁶, M. J. Da Cunha Sargedas De Sousa^{60b}, J. V. Da Fonseca Pinto^{81b}, C. Da Via¹⁰¹, W. Dabrowski^{84a}, T. Dado^{28a}, S. Dahbi^{35e}, T. Dai¹⁰⁶, C. Dallapiccola¹⁰³, M. Dam⁴⁰, G. D'amen²⁹, V. D'Amico^{75a,75b}, J. Damp¹⁰⁰, J. R. Dandoy¹³⁷, M. F. Daneri³⁰, N. P. Dang^{181,i}, N. S. Dann¹⁰¹, M. Danninger¹⁷⁵, V. Dao³⁶, G. Darbo^{55b}, O. Dartsis⁵, A. Dattagupta¹³², T. Daubney⁴⁶, S. D'Auria^{69a,69b}, W. Davey²⁴, C. David⁴⁶, T. Davidek¹⁴³, D. R. Davis⁴⁹, I. Dawson¹⁴⁹, K. De⁸, R. De Asmundis^{70a}, M. De Beurs¹²⁰, S. De Castro^{23a,23b}, S. De Cecco^{73a,73b}, N. De Groot¹¹⁹, P. de Jong¹²⁰, H. De la Torre¹⁰⁷, A. De Maria^{15c}, D. De Pedis^{73a}, A. De Salvo^{73a}, U. De Sanctis^{74a,74b}, M. De Santis^{74a,74b}, A. De Santo¹⁵⁶, K. De Vasconcelos Corga¹⁰², J. B. De Vivie De Regie⁶⁵, C. Debenedetti¹⁴⁶, D. V. Dedovich⁸⁰, A. M. Deiana⁴², M. Del Gaudio^{41a,41b}, J. Del Peso⁹⁹, Y. Delabat Diaz⁴⁶, D. Delgove⁶⁵, F. Deliot^{145,q}, C. M. Delitzsch⁷, M. Della Pietra^{70a,70b}, D. Della Volpe⁵⁴, A. Dell'Acqua³⁶, L. Dell'Asta^{74a,74b}, M. Delmastro⁵, C. Delporte⁶⁵, P. A. Delsart⁵⁸, D. A. DeMarco¹⁶⁷, S. Demers¹⁸³, M. Demichev⁸⁰, G. Demontigny¹¹⁰, S. P. Denisov¹²³, D. Denysiuk¹²⁰, L. D'Eramo¹³⁶, D. Derendarz⁸⁵, J. E. Derkaoui^{35d}, F. Derue¹³⁶, P. Dervan⁹¹, K. Desch²⁴, C. Deterre⁴⁶, K. Dette¹⁶⁷, C. Deutsch²⁴, M. R. Devesa³⁰, P. O. Deviveiros³⁶, A. Dewhurst¹⁴⁴, F. A. Di Bello⁵⁴, A. Di Ciaccio^{74a,74b}, L. Di Ciaccio⁵, W. K. Di Clemente¹³⁷, C. Di Donato^{70a,70b}, A. Di Girolamo³⁶, G. Di Gregorio^{72a,72b}, B. Di Micco^{75a,75b}, R. Di Nardo¹⁰³, K. F. Di Petrillo⁵⁹, R. Di Sipio¹⁶⁷, D. Di Valentino³⁴, C. Diaconu¹⁰², F. A. Dias⁴⁰, T. Dias Do Vale^{140a}, M. A. Diaz^{147a}, J. Dickinson¹⁸, E. B. Diehl¹⁰⁶, J. Dietrich¹⁹, S. Díez Cornell⁴⁶, A. Dimitrievska¹⁸, W. Ding^{15b}, J. Dingfelder²⁴, F. Dittus³⁶, F. Djama¹⁰², T. Djobava^{159b}, J. I. Djuvsland¹⁷, M. A. B. Do Vale^{81c}, M. Dobre^{27b}, D. Dodsworth²⁶, C. Doglioni⁹⁷, J. Dolejsi¹⁴³, Z. Dolezal¹⁴³, M. Donadelli^{81d}, B. Dong^{60c}, J. Donini³⁸, A. D'onofrio⁹³, M. D'Onofrio⁹¹, J. Dopke¹⁴⁴, A. Doria^{70a}, M. T. Dova⁸⁹, A. T. Doyle⁵⁷, E. Drechsler¹⁵², E. Dreyer¹⁵², T. Dreyer⁵³, A. S. Drobac¹⁷⁰, D. Du^{60b}, Y. Duan^{60b}, F. Dubinin¹¹¹, M. Dubovsky^{28a}, A. Dubreuil⁵⁴, E. Duchovni¹⁸⁰, G. Duckeck¹¹⁴, A. Ducourthial¹³⁶, O. A. Ducu¹¹⁰, D. Duda¹¹⁵, A. Dudarev³⁶, A. C. Dudder¹⁰⁰, E. M. Duffield¹⁸, L. Duflot⁶⁵, M. Dührssen³⁶, C. Dülsen¹⁸², M. Dumancic¹⁸⁰, A. E. Dumitriu^{27b}, A. K. Duncan⁵⁷, M. Dunford^{61a}, A. Duperrin¹⁰², H. Duran Yildiz^{4a}, M. Düren⁵⁶, A. Durglishvili^{159b}, D. Duschinger⁴⁸, B. Dutta⁴⁶, D. Duvnjak¹, G. I. Dyckes¹³⁷, M. Dyndal³⁶, S. Dysch¹⁰¹, B. S. Dziedzic⁸⁵, K. M. Ecker¹¹⁵, R. C. Edgar¹⁰⁶, M. G. Eggleston⁴⁹, T. Eifert³⁶, G. Eigen¹⁷, K. Einsweiler¹⁸, T. Ekelof¹⁷², H. El Jarrari^{35e}, M. El Kacimi^{35c}, R. El Kosseifi¹⁰², V. Ellajosyula¹⁷², M. Ellert¹⁷², F. Ellinghaus¹⁸², A. A. Elliot⁹³, N. Ellis³⁶, J. Elmsheuser²⁹, M. Elsing³⁶, D. Emeliyanov¹⁴⁴, A. Emerman³⁹, Y. Enari¹⁶³, M. B. Epland⁴⁹, J. Erdmann⁴⁷, A. Ereditato²⁰, M. Errenst³⁶, M. Escalier⁶⁵, C. Escobar¹⁷⁴, O. Estrada Pastor¹⁷⁴, E. Etzion¹⁶¹, H. Evans⁶⁶, A. Ezhilov¹³⁸, F. Fabbri⁵⁷, L. Fabbri^{23a,23b}, V. Fabiani¹¹⁹, G. Facini⁹⁵, R. M. Faisca Rodrigues Pereira^{140a}, R. M. Fakhrutdinov¹²³, S. Falciano^{73a}, P. J. Falke⁵, S. Falke⁵, J. Faltova¹⁴³, Y. Fang^{15a}, Y. Fang^{15a}, G. Fanourakis⁴⁴, M. Fanti^{69a,69b}, M. Faraj^{67a,67c,t}, A. Farbin⁸, A. Farilla^{75a}, E. M. Farina^{71a,71b}, T. Farooque¹⁰⁷, S. Farrell¹⁸, S. M. Farrington⁵⁰, P. Farthouat³⁶, F. Fassi^{35e}, P. Fassnacht³⁶, D. Fassouliotis⁹, M. Fauci Giannelli⁵⁰, W. J. Fawcett³², L. Fayard⁶⁵, O. L. Fedin^{138,o}, W. Fedorko¹⁷⁵, A. Fehr²⁰, M. Feickert⁴², L. Felgioni¹⁰², A. Fell¹⁴⁹, C. Feng^{60b}, E. J. Feng³⁶, M. Feng⁴⁹, M. J. Fenton⁵⁷, A. B. Fenyuk¹²³, J. Ferrando⁴⁶, A. Ferrante¹⁷³, A. Ferrari¹⁷², P. Ferrari¹²⁰, R. Ferrari^{71a}, D. E. Ferreira de Lima^{61b}, A. Ferrer¹⁷⁴, D. Ferrere⁵⁴, C. Ferretti¹⁰⁶, F. Fiedler¹⁰⁰, A. Filipčić⁹², F. Filthaut¹¹⁹, K. D. Finelli²⁵, M. C. N. Fiolhais^{140a,140c,a}, L. Fiorini¹⁷⁴, F. Fischer¹¹⁴, W. C. Fisher¹⁰⁷, I. Fleck¹⁵¹, P. Fleischmann¹⁰⁶, R. R. M. Fletcher¹³⁷, T. Flick¹⁸², B. M. Flierl¹¹⁴, L. Flores¹³⁷, L. R. Flores Castillo^{63a}, F. M. Follega^{76a,76b}, N. Fomin¹⁷, J. H. Foo¹⁶⁷, G. T. Forcolin^{76a,76b}, A. Formica¹⁴⁵, F. A. Förster¹⁴, A. C. Forti¹⁰¹, A. G. Foster²¹, M. G. Foti¹³⁵, D. Fournier⁶⁵, H. Fox⁹⁰, P. Francavilla^{72a,72b}, S. Francescato^{73a,73b}, M. Franchini^{23a,23b}, S. Franchino^{61a}, D. Francis³⁶, L. Franconi²⁰, M. Franklin⁵⁹,

A. N. Fray⁹³, P. M. Freeman²¹, B. Freund¹¹⁰, W. S. Freund^{81b}, E. M. Freundlich⁴⁷, D. C. Frizzell¹²⁹, D. Froidevaux³⁶, J. A. Frost¹³⁵, C. Fukunaga¹⁶⁴, E. Fullana Torregrosa¹⁷⁴, E. Fumagalli^{55a,55b}, T. Fusayasu¹¹⁶, J. Fuster¹⁷⁴, A. Gabrielli^{23a,23b}, A. Gabrielli¹⁸, G. P. Gach^{84a}, S. Gadatsch⁵⁴, P. Gadow¹¹⁵, G. Gagliardi^{55a,55b}, L. G. Gagnon¹¹⁰, C. Galea^{27b}, B. Galhardo^{140a}, G. E. Gallardo¹³⁵, E. J. Gallas¹³⁵, B. J. Gallop¹⁴⁴, G. Galster⁴⁰, R. Gamboa Goni⁹³, K. K. Gan¹²⁷, S. Ganguly¹⁸⁰, J. Gao^{60a}, Y. Gao⁵⁰, Y. S. Gao^{31.l}, C. García¹⁷⁴, J. E. García Navarro¹⁷⁴, J. A. García Pascual^{15a}, C. Garcia-Argos⁵², M. Garcia-Sciveres¹⁸, R. W. Gardner³⁷, N. Garelli¹⁵³, S. Gargiulo⁵², V. Garonne¹³⁴, P. Gaspar^{81b}, A. Gaudiello^{55a,55b}, G. Gaudio^{71a}, I. L. Gavrilenko¹¹¹, A. Gavrilyuk¹²⁴, C. Gay¹⁷⁵, G. Gaycken⁴⁶, E. N. Gazis¹⁰, A. A. Geanta^{27b}, C. M. Gee¹⁴⁶, C. N. P. Gee¹⁴⁴, J. Geisen⁵³, M. Geisen¹⁰⁰, C. Gemme^{55b}, M. H. Genest⁵⁸, C. Geng¹⁰⁶, S. Gentile^{73a,73b}, S. George⁹⁴, T. Gerialis⁴⁴, L. O. Gerlach⁵³, P. Gessinger-Befurt¹⁰⁰, G. Gessner⁴⁷, S. Ghasemi¹⁵¹, M. Ghasemi Bostanabad¹⁷⁶, M. Ghneimat¹⁵¹, A. Ghosh⁶⁵, A. Ghosh⁷⁸, B. Giacobbe^{23b}, S. Giagu^{73a,73b}, N. Giangiacomi^{23a,23b}, P. Giannetti^{72a}, A. Giannini^{70a,70b}, G. Giannini¹⁴, S. M. Gibson⁹⁴, M. Gignac¹⁴⁶, D. Gillberg³⁴, G. Gilles¹⁸², D. M. Gingrich^{3,au}, M. P. Giordani^{67a,67c}, F. M. Giorgi^{23b}, P. F. Giraud¹⁴⁵, G. Giugliarelli^{67a,67c}, D. Giugni^{69a}, F. Giulì^{74a,74b}, S. Gkaitatzis¹⁶², I. Gkialas^{9.g}, E. L. Gkoukousis¹⁴, P. Gkoutoumis¹⁰, L. K. Gladilin¹¹³, C. Glasman⁹⁹, J. Glatzer¹⁴, P. C. F. Glaysheer⁴⁶, A. Glazov⁴⁶, G. R. Gledhill¹³², M. Goblirsch-Kolb²⁶, D. Godin¹¹⁰, S. Goldfarb¹⁰⁵, T. Golling⁵⁴, D. Golubkov¹²³, A. Gomes^{140a,140b}, R. Goncalves Gama⁵³, R. Gonçalo^{140a}, G. Gonella⁵², L. Gonella²¹, A. Gongadze⁸⁰, F. Gonnella²¹, J. L. Gonski⁵⁹, S. González de la Hoz¹⁷⁴, S. Gonzalez-Sevilla⁵⁴, G. R. Gonzalvo Rodriguez¹⁷⁴, L. Goossens³⁶, P. A. Gorbounov¹²⁴, H. A. Gordon²⁹, B. Gorini³⁶, E. Gorini^{68a,68b}, A. Gorišek⁹², A. T. Goshaw⁴⁹, M. I. Gostkin⁸⁰, C. A. Gottardo¹¹⁹, M. Goughri^{35b}, D. Goujdami^{35c}, A. G. Goussiou¹⁴⁸, N. Govender^{33c}, C. Goy⁵, E. Gozani¹⁶⁰, I. Grabowska-Bold^{84a}, E. C. Graham⁹¹, J. Gramling¹⁷¹, E. Gramstad¹³⁴, S. Grancagnolo¹⁹, M. Grandi¹⁵⁶, V. Gratchev¹³⁸, P. M. Gravila^{27f}, F. G. Gravili^{68a,68b}, C. Gray⁵⁷, H. M. Gray¹⁸, C. Grefe²⁴, K. Gregersen⁹⁷, I. M. Gregor⁴⁶, P. Grenier¹⁵³, K. Grevtsov⁴⁶, C. Grieco¹⁴, N. A. Grieser¹²⁹, A. A. Grillo¹⁴⁶, K. Grimm^{31.k}, S. Grinstein^{14.z}, J.-F. Grivaz⁶⁵, S. Groh¹⁰⁰, E. Gross¹⁸⁰, J. Grosse-Knetter⁵³, Z. J. Grout⁹⁵, C. Grud¹⁰⁶, A. Grummer¹¹⁸, L. Guan¹⁰⁶, W. Guan¹⁸¹, C. Gubbels¹⁷⁵, J. Guenther³⁶, A. Guerguichon⁶⁵, J. G. R. Guerrero Rojas¹⁷⁴, F. Guescini¹¹⁵, D. Guest¹⁷¹, R. Gugel⁵², T. Guillemain⁵, S. Guindon³⁶, U. Gul⁵⁷, J. Guo^{60c}, W. Guo¹⁰⁶, Y. Guo^{60a.s}, Z. Guo¹⁰², R. Gupta⁴⁶, S. Gurbuz^{12c}, G. Gustavino¹²⁹, M. Guth⁵², P. Gutierrez¹²⁹, C. Gutschow⁹⁵, C. Guyot¹⁴⁵, C. Gwenlan¹³⁵, C. B. Gwilliam⁹¹, A. Haas¹²⁵, C. Haber¹⁸, H. K. Hadavand⁸, N. Haddad^{35e}, A. Hadeef^{60a}, S. Hageböck³⁶, M. Haleem¹⁷⁷, J. Haley¹³⁰, G. Halladjian¹⁰⁷, G. D. Hallowell¹⁰², K. Hamacher¹⁸², P. Hamal¹³¹, K. Hamano¹⁷⁶, H. Hamdaoui^{35e}, M. Hamer²⁴, G. N. Hamity¹⁴⁹, K. Han^{60a.y}, L. Han^{60a}, S. Han^{15a}, Y. F. Han¹⁶⁷, K. Hanagaki^{82.w}, M. Hance¹⁴⁶, D. M. Handl¹¹⁴, B. Haney¹³⁷, R. Hankache¹³⁶, E. Hansen⁹⁷, J. B. Hansen⁴⁰, J. D. Hansen⁴⁰, M. C. Hansen²⁴, P. H. Hansen⁴⁰, E. C. Hanson¹⁰¹, K. Hara¹⁶⁹, T. Harenberg¹⁸², S. Harkusha¹⁰⁸, P. F. Harrison¹⁷⁸, N. M. Hartmann¹¹⁴, Y. Hasegawa¹⁵⁰, A. Hasib⁵⁰, S. Hassani¹⁴⁵, S. Haug²⁰, R. Hauser¹⁰⁷, L. B. Havener³⁹, M. Havranek¹⁴², C. M. Hawkes²¹, R. J. Hawkins³⁶, D. Hayden¹⁰⁷, C. Hayes¹⁵⁵, R. L. Hayes¹⁷⁵, C. P. Hays¹³⁵, J. M. Hays⁹³, H. S. Hayward⁹¹, S. J. Haywood¹⁴⁴, F. He^{60a}, M. P. Heath⁵⁰, V. Hedberg⁹⁷, L. Heelan⁸, S. Heer²⁴, K. K. Heidegger⁵², W. D. Heidorn⁷⁹, J. Heilman³⁴, S. Heim⁴⁶, T. Heim¹⁸, B. Heinemann^{46.ar}, J. J. Heinrich¹³², L. Heinrich³⁶, C. Heinz⁵⁶, J. Hejbal¹⁴¹, L. Helary^{61b}, A. Held¹⁷⁵, S. Hellesund¹³⁴, C. M. Helling¹⁴⁶, S. Hellman^{45a,45b}, C. Helsens³⁶, R. C. W. Henderson⁹⁰, Y. Heng¹⁸¹, S. Henkelmann¹⁷⁵, A. M. Henriques Correia³⁶, G. H. Herbert¹⁹, H. Herde²⁶, V. Herget¹⁷⁷, Y. Hernández Jiménez^{33e}, H. Herr¹⁰⁰, M. G. Herrmann¹¹⁴, T. Herrmann⁴⁸, G. Herten⁵², R. Hertenberger¹¹⁴, L. Hervas³⁶, T. C. Herwig¹³⁷, G. G. Hesketh⁹⁵, N. P. Hessey^{168a}, A. Higashida¹⁶³, S. Higashino⁸², E. Higón-Rodríguez¹⁷⁴, K. Hildebrand³⁷, E. Hill¹⁷⁶, J. C. Hill³², K. K. Hill²⁹, K. H. Hiller⁴⁶, S. J. Hillier²¹, M. Hils⁴⁸, I. Hinchliffe¹⁸, F. Hinterkeuser²⁴, M. Hirose¹³³, S. Hirose⁵², D. Hirschbuehl¹⁸², B. Hiti⁹², O. Hladik¹⁴¹, D. R. Hlalu^{33e}, X. Hoad⁵⁰, J. Hobbs¹⁵⁵, N. Hod¹⁸⁰, M. C. Hodgkinson¹⁴⁹, A. Hoecker³⁶, F. Hoenig¹¹⁴, D. Hohn⁵², D. Hohov⁶⁵, T. Holm²⁴, T. R. Holmes³⁷, M. Holzbock¹¹⁴, L. B. A. H. Hommels³², S. Honda¹⁶⁹, T. M. Hong¹³⁹, J. C. Honig⁵², A. Hönle¹¹⁵, B. H. Hoerberman¹⁷³, W. H. Hopkins⁶, Y. Horii¹¹⁷, P. Horn⁴⁸, L. A. Horyn³⁷, S. Hou¹⁵⁸, A. Houmada^{35a}, J. Howarth¹⁰¹, J. Hoya⁸⁹, M. Hrabovsky¹³¹, J. Hrdinka⁷⁷, I. Hristova¹⁹, J. Hrivnac⁶⁵, A. Hrynevich¹⁰⁹, T. Hryn'ova⁵, P. J. Hsu⁶⁴, S.-C. Hsu¹⁴⁸, Q. Hu²⁹, S. Hu^{60c}, Y. F. Hu^{15a,15d}, D. P. Huang⁹⁵, Y. Huang^{60a}, Y. Huang^{15a}, Z. Hubacek¹⁴², F. Hubaut¹⁰², M. Huebner²⁴, F. Huegging²⁴, T. B. Huffman¹³⁵, M. Huhtinen³⁶, R. F. H. Hunter³⁴, P. Huo¹⁵⁵, A. M. Hupe³⁴, N. Huseynov^{80.ag}, J. Huston¹⁰⁷, J. Huth⁵⁹, R. Hyneman¹⁰⁶, S. Hyrych^{28a}, G. Iacobucci⁵⁴, G. Iakovidis²⁹, I. Ibragimov¹⁵¹, L. Iconomidou-Fayard⁶⁵, Z. Idrissi^{35e},


P. Iengo³⁶, R. Ignazzi⁴⁰, O. Igonkina^{120,ab}, R. Iguchi¹⁶³, T. Iizawa⁵⁴, Y. Ikegami⁸², M. Ikeno⁸², D. Iliadis¹⁶², N. Ilic^{119,167,ae}, F. Iltzsche⁴⁸, G. Introzzi^{71a,71b}, M. Iodice^{75a}, K. Iordanidou^{168a}, V. Ippolito^{73a,73b}, M. F. Isacson¹⁷², M. Ishino¹⁶³, W. Islam¹³⁰, C. Issever^{19,46}, S. Istin¹⁶⁰, F. Ito¹⁶⁹, J. M. Iturbe Ponce^{63a}, R. Iuppa^{76a,76b}, A. Ivina¹⁸⁰, H. Iwasaki⁸², J. M. Izen⁴³, V. Izzo^{70a}, P. Jacka¹⁴¹, P. Jackson¹, R. M. Jacobs²⁴, B. P. Jaeger¹⁵², V. Jain², G. Jäkel¹⁸², K. B. Jakobi¹⁰⁰, K. Jakobs⁵², T. Jakoubek¹⁴¹, J. Jamieson⁵⁷, K. W. Janas^{84a}, R. Jansky⁵⁴, J. Janssen²⁴, M. Janus⁵³, P. A. Janus^{84a}, G. Jarlskog⁹⁷, N. Javadov^{80,ag}, T. Javůrek³⁶, M. Javurkova⁵², F. Jeanneau¹⁴⁵, L. Jeanty¹³², J. Jejelava^{159a,ah}, A. Jelinskas¹⁷⁸, P. Jenni^{52,b}, J. Jeong⁴⁶, N. Jeong⁴⁶, S. Jézéquel⁵, H. Ji¹⁸¹, J. Jia¹⁵⁵, H. Jiang⁷⁹, Y. Jiang^{60a}, Z. Jiang^{153,p}, S. Jiggins⁵², F. A. Jimenez Morales³⁸, J. Jimenez Pena¹¹⁵, S. Jin^{15c}, A. Jinaru^{27b}, O. Jinnouchi¹⁶⁵, H. Jivan^{33e}, P. Johansson¹⁴⁹, K. A. Johns⁷, C. A. Johnson⁶⁶, K. Jon-And^{45a,45b}, R. W. L. Jones⁹⁰, S. D. Jones¹⁵⁶, S. Jones⁷, T. J. Jones⁹¹, J. Jongmanns^{61a}, P. M. Jorge^{140a}, J. Jovicevic³⁶, X. Ju¹⁸, J. J. Junggeburth¹¹⁵, A. Juste Rozas^{14,z}, A. Kaczmarek⁸⁵, M. Kado^{73a,73b}, H. Kagan¹²⁷, M. Kagan¹⁵³, A. Kahn³⁹, C. Kahra¹⁰⁰, T. Kaji¹⁷⁹, E. Kajomovitz¹⁶⁰, C. W. Kalderon⁹⁷, A. Kaluza¹⁰⁰, A. Kamenshchikov¹²³, M. Kaneda¹⁶³, N. J. Kang¹⁴⁶, L. Kanjir⁹², Y. Kano¹¹⁷, V. A. Kantserov¹¹², J. Kanzaki⁸², L. S. Kaplan¹⁸¹, D. Kar^{33e}, K. Karava¹³⁵, M. J. Kareem^{168b}, S. N. Karpov⁸⁰, Z. M. Karpova⁸⁰, V. Kartvelishvili⁹⁰, A. N. Karyukhin¹²³, L. Kashif¹⁸¹, R. D. Kass¹²⁷, A. Kastanas^{45a,45b}, C. Kato^{60c,60d}, J. Katzy⁴⁶, K. Kawade¹⁵⁰, K. Kawagoe⁸⁸, T. Kawaguchi¹¹⁷, T. Kawamoto¹⁶³, G. Kawamura⁵³, E. F. Kay¹⁷⁶, V. F. Kazanin^{122a,122b}, R. Keeler¹⁷⁶, R. Kehoe⁴², J. S. Keller³⁴, E. Kellermann⁹⁷, D. Kelsey¹⁵⁶, J. J. Kempster²¹, J. Kendrick²¹, K. E. Kennedy³⁹, O. Kepka¹⁴¹, S. Kersten¹⁸², B. P. Kerševan⁹², S. Ketabchi Haghighat¹⁶⁷, M. Khader¹⁷³, F. Khalil-Zada¹³, M. Khandoga¹⁴⁵, A. Khanov¹³⁰, A. G. Kharlamov^{122b,122a}, T. Kharlamova^{122b,122a}, E. E. Khoda¹⁷⁵, A. Khodinov¹⁶⁶, T. J. Khoo⁵⁴, E. Khranov⁸⁰, J. Khubua^{159b}, S. Kido⁸³, M. Kiehn⁵⁴, C. R. Kilby⁹⁴, Y. K. Kim³⁷, N. Kimura⁹⁵, O. M. Kind¹⁹, B. T. King⁹¹, D. Kirchmeier⁴⁸, J. Kirk¹⁴⁴, A. E. Kiryunin¹¹⁵, T. Kishimoto¹⁶³, D. P. Kisiuk¹⁶⁷, V. Kitali⁴⁶, O. Kiverny⁵, T. Klapdor-Kleingrothaus⁵², M. Klassen^{61a}, M. H. Klein¹⁰⁶, M. Klein⁹¹, U. Klein⁹¹, K. Kleinknecht¹⁰⁰, P. Klimek¹²¹, A. Klimentov²⁹, T. Klingl²⁴, T. Klioutchnikova³⁶, F. F. Klitzner¹¹⁴, P. Kluit¹²⁰, S. Kluth¹¹⁵, E. Kneringer⁷⁷, E. B. F. G. Knoop¹⁰², A. Knue⁵², D. Kobayashi⁸⁸, T. Kobayashi¹⁶³, M. Kobel⁴⁸, M. Kocian¹⁵³, P. Kodys¹⁴³, P. T. Koenig²⁴, T. Koffas³⁴, N. M. Köhler³⁶, T. Koi¹⁵³, M. Kolb^{61b}, I. Koletsou⁵, T. Komarek¹³¹, T. Kondo⁸², N. Kondrashova^{60c}, K. Köneke⁵², A. C. König¹¹⁹, T. Kono¹²⁶, R. Konoplich^{125,am}, V. Konstantinides⁹⁵, N. Konstantinidis⁹⁵, B. Konya⁹⁷, R. Kopeliansky⁶⁶, S. Koperny^{84a}, K. Korcyl⁸⁵, K. Kordas¹⁶², G. Koren¹⁶¹, A. Korn⁹⁵, I. Korolkov¹⁴, E. V. Korolkova¹⁴⁹, N. Korotkova¹¹³, O. Kortner¹¹⁵, S. Kortner¹¹⁵, T. Kosek¹⁴³, V. V. Kostyukhin¹⁶⁶, A. Kotskechagia⁶⁵, A. Kotwal⁴⁹, A. Koulouris¹⁰, A. Kourkoumeli-Charalampidi^{71a,71b}, C. Kourkoumelis⁹, E. Kourlitis¹⁴⁹, V. Kouskoura²⁹, A. B. Kowalewska⁸⁵, R. Kowalewski¹⁷⁶, C. Kozakai¹⁶³, W. Kozanecki¹⁴⁵, A. S. Kozhin¹²³, V. A. Kramarenko¹¹³, G. Kramberger⁹², D. Krasnopevtsev^{60a}, M. W. Krasny¹³⁶, A. Krasznahorkay³⁶, D. Krauss¹¹⁵, J. A. Kremer^{84a}, J. Kretschmar⁹¹, P. Krieger¹⁶⁷, F. Krieter¹¹⁴, A. Krishnan^{61b}, K. Krizka¹⁸, K. Kroeninger⁴⁷, H. Kroha¹¹⁵, J. Kroll¹⁴¹, J. Kroll¹³⁷, K. S. Krowpman¹⁰⁷, J. Krstic¹⁶, U. Kruchonak⁸⁰, H. Krüger²⁴, N. Krumnack⁷⁹, M. C. Kruse⁴⁹, J. A. Krzysiak⁸⁵, T. Kubota¹⁰⁵, O. Kuchinskaia¹⁶⁶, S. Kuday^{4b}, J. T. Kuechler⁴⁶, S. Kuehn³⁶, A. Kugel^{61a}, T. Kuhl⁴⁶, V. Kukhtin⁸⁰, R. Kukla¹⁰², Y. Kulchitsky^{108,aj}, S. Kuleshov^{147d}, Y. P. Kulinich¹⁷³, M. Kuna⁵⁸, T. Kunigo⁸⁶, A. Kupco¹⁴¹, T. Kupfer⁴⁷, O. Kuprash⁵², H. Kurashige⁸³, L. L. Kurchaninov^{168a}, Y. A. Kurochkin¹⁰⁸, A. Kurova¹¹², M. G. Kurth^{15a,15d}, E. S. Kuwertz³⁶, M. Kuze¹⁶⁵, A. K. Kvam¹⁴⁸, J. Kvita¹³¹, T. Kwan¹⁰⁴, A. La Rosa¹¹⁵, L. La Rotonda^{41b,41a}, F. La Ruffa^{41a,41b}, C. Lacasta¹⁷⁴, F. Lacava^{73a,73b}, D. P. J. Lack¹⁰¹, H. Lacker¹⁹, D. Lacour¹³⁶, E. Ladygin⁸⁰, R. Lafaye⁵, B. Laforge¹³⁶, T. Lagouri^{33e}, S. Lai⁵³, S. Lammers⁶⁶, W. Lampl⁷, C. Lampoudis¹⁶², E. Lançon²⁹, U. Landgraf⁵², M. P. J. Landon⁹³, M. C. Lanfermann⁵⁴, V. S. Lang⁴⁶, J. C. Lange⁵³, R. J. Langenberg³⁶, A. J. Lankford¹⁷¹, F. Lanni²⁹, K. Lantzsch²⁴, A. Lanza^{71a}, A. Lapertosa^{55a,55b}, S. Laplace¹³⁶, J. F. Laporte¹⁴⁵, T. Lari^{69a}, F. Lasagni Manghi^{23a,23b}, M. Lassnig³⁶, T. S. Lau^{63a}, A. Laudrain⁶⁵, A. Laurier³⁴, M. Lavorgna^{70a,70b}, S. D. Lawlor⁹⁴, M. Lazzaroni^{69a,69b}, B. Le¹⁰⁵, E. Le Guirriec¹⁰², M. LeBlanc⁷, T. LeCompte⁶, F. Ledroit-Guillon⁵⁸, A. C. A. Lee⁹⁵, C. A. Lee²⁹, G. R. Lee¹⁷, L. Lee⁵⁹, S. C. Lee¹⁵⁸, S. J. Lee³⁴, S. Lee⁷⁹, B. Lefebvre^{168a}, H. P. Lefebvre⁹⁴, M. Lefebvre¹⁷⁶, F. Legger¹¹⁴, C. Leggett¹⁸, K. Lehmann¹⁵², N. Lehmann¹⁸², G. Lehmann Miotto³⁶, W. A. Leight⁴⁶, A. Leisos^{162,x}, M. A. L. Leite^{81d}, C. E. Leitgeb¹¹⁴, R. Leitner¹⁴³, D. Lellouch¹⁸⁰, K. J. C. Leney⁴², T. Lenz²⁴, B. Lenzi³⁶, R. Leone⁷, S. Leone^{72a}, C. Leonidopoulos⁵⁰, A. Leopold¹³⁶, G. Lerner¹⁵⁶, C. Leroy¹¹⁰, R. Les¹⁶⁷, C. G. Lester³², M. Levchenko¹³⁸, J. Levêque⁵, D. Levin¹⁰⁶, L. J. Levinson¹⁸⁰, D. J. Lewis²¹, B. Li^{15b}, B. Li¹⁰⁶, C.-Q. Li^{60a}, F. Li^{60c}, H. Li^{60a}, H. Li^{60b}, J. Li^{60c}

K. Li¹⁵³, L. Li^{60c}, M. Li^{15a,15d}, Q. Li^{15a,15d}, Q. Y. Li^{60a}, S. Li^{60c,60d}, X. Li⁴⁶, Y. Li⁴⁶, Z. Li^{60b}, Z. Liang^{15a}, B. Liberti^{74a}, A. Liblong¹⁶⁷, K. Lie^{63c}, S. Lim²⁹, C. Y. Lin³², K. Lin¹⁰⁷, T. H. Lin¹⁰⁰, R. A. Linck⁶⁶, J. H. Lindon²¹, A. L. Lioni⁵⁴, E. Lipeles¹³⁷, A. Lipniacka¹⁷, M. Lisovyi^{61b}, T. M. Liss^{173.at}, A. Lister¹⁷⁵, A. M. Litke¹⁴⁶, J. D. Little⁸, B. Liu⁷⁹, B. L. Liu⁶, H. B. Liu²⁹, H. Liu¹⁰⁶, J. B. Liu^{60a}, J. K. K. Liu¹³⁵, K. Liu¹³⁶, M. Liu^{60a}, P. Liu¹⁸, Y. Liu^{15a,15d}, Y. L. Liu¹⁰⁶, Y. W. Liu^{60a}, M. Livan^{71a,71b}, A. Lleres⁵⁸, J. Llorente Merino¹⁵², S. L. Lloyd⁹³, C. Y. Lo^{63b}, F. Lo Sterzo⁴², E. M. Lobodzinska⁴⁶, P. Loch⁷, S. Loffredo^{74a,74b}, T. Lohse¹⁹, K. Lohwasser¹⁴⁹, M. Lokajicek¹⁴¹, J. D. Long¹⁷³, R. E. Long⁹⁰, L. Longo³⁶, K. A. Looper¹²⁷, J. A. Lopez^{147d}, I. Lopez Paz¹⁰¹, A. Lopez Solis¹⁴⁹, J. Lorenz¹¹⁴, N. Lorenzo Martinez⁵, M. Losada^{22a}, P. J. Lösel¹¹⁴, A. Lösle⁵², X. Lou⁴⁶, X. Lou^{15a}, A. Lounis⁶⁵, J. Love⁶, P. A. Love⁹⁰, J. J. Lozano Bahilo¹⁷⁴, M. Lu^{60a}, Y. J. Lu⁶⁴, H. J. Lubatti¹⁴⁸, C. Luci^{73a,73b}, A. Lucotte⁵⁸, C. Luedtke⁵², F. Luehring⁶⁶, I. Luise¹³⁶, L. Luminari^{73a}, B. Lund-Jensen¹⁵⁴, M. S. Lutz¹⁰³, D. Lynn²⁹, H. Lyons⁹¹, R. Lysak¹⁴¹, E. Lytken⁹⁷, F. Lyu^{15a}, V. Lyubushkin⁸⁰, T. Lyubushkina⁸⁰, H. Ma²⁹, L. L. Ma^{60b}, Y. Ma^{60b}, G. Maccarrone⁵¹, A. Macchiolo¹¹⁵, C. M. Macdonald¹⁴⁹, J. Machado Miguens¹³⁷, D. Madaffari¹⁷⁴, R. Madar³⁸, W. F. Mader⁴⁸, N. Madysa⁴⁸, J. Maeda⁸³, T. Maeno²⁹, M. Maerker⁴⁸, A. S. Maevskiy¹¹³, V. Magerl⁵², N. Magini⁷⁹, D. J. Mahon³⁹, C. Maidantchik^{81b}, T. Maier¹¹⁴, A. Maio^{140a,140b,140d}, K. Maj^{84a}, O. Majersky^{28a}, S. Majewski¹³², Y. Makida⁸², N. Makovec⁶⁵, B. Malaescu¹³⁶, Pa. Malecki⁸⁵, V. P. Maleev¹³⁸, F. Malek⁵⁸, U. Mallik⁷⁸, D. Malon⁶, C. Malone³², S. Maltezos¹⁰, S. Malyukov⁸⁰, J. Mamuzic¹⁷⁴, G. Mancini⁵¹, I. Mandić⁹², L. Manhaes de Andrade Filho^{81a}, I. M. Maniatis¹⁶², J. Manjarres Ramos⁴⁸, K. H. Mankinen⁹⁷, A. Mann¹¹⁴, A. Manousos⁷⁷, B. Mansoulie¹⁴⁵, I. Manthos¹⁶², S. Manzoni¹²⁰, A. Marantis¹⁶², G. Marceca³⁰, L. Marchese¹³⁵, G. Marchiori¹³⁶, M. Marcisovsky¹⁴¹, L. Marcocchia^{74a,74b}, C. Marcon⁹⁷, C. A. Marin Tobon³⁶, M. Marjanovic¹²⁹, Z. Marshall¹⁸, M. U. F. Martensson¹⁷², S. Marti-Garcia¹⁷⁴, C. B. Martin¹²⁷, T. A. Martin¹⁷⁸, V. J. Martin⁵⁰, B. Martin dit Latour¹⁷, L. Martinelli^{75a,75b}, M. Martinez^{14.z}, V. I. Martinez Outschoorn¹⁰³, S. Martin-Haugh¹⁴⁴, V. S. Martoiu^{27b}, A. C. Martyniuk⁹⁵, A. Marzin³⁶, S. R. Maschek¹¹⁵, L. Masetti¹⁰⁰, T. Mashimo¹⁶³, R. Mashinistov¹¹¹, J. Masik¹⁰¹, A. L. Maslennikov^{122b,122a}, L. Massa^{74a,74b}, P. Massarotti^{70a,70b}, P. Mastrandrea^{72a,72b}, A. Mastroberardino^{41a,41b}, T. Masubuchi¹⁶³, D. Matakias¹⁰, A. Matic¹¹⁴, N. Matsuzawa¹⁶³, P. Mättig²⁴, J. Maurer^{27b}, B. Maček⁹², D. A. Maximov^{122a,122b}, R. Mazini¹⁵⁸, I. Maznas¹⁶², S. M. Mazza¹⁴⁶, S. P. Mc Kee¹⁰⁶, T. G. McCarthy¹¹⁵, W. P. McCormack¹⁸, E. F. McDonald¹⁰⁵, J. A. McFayden³⁶, G. Mchedlidze^{159b}, M. A. McKay⁴², K. D. McLean¹⁷⁶, S. J. McMahon¹⁴⁴, P. C. McNamara¹⁰⁵, C. J. McNicol¹⁷⁸, R. A. McPherson^{176.ac}, J. E. Mdhului^{33e}, Z. A. Meadows¹⁰³, S. Meehan³⁶, T. Megy⁵², S. Mehlhase¹¹⁴, A. Mehta⁹¹, T. Meideck⁵⁸, B. Meirose⁴³, D. Melini¹⁷⁴, B. R. Mellado Garcia^{33e}, J. D. Mellenthin⁵³, M. Melo^{28a}, F. Meloni⁴⁶, A. Melzer²⁴, S. B. Menary¹⁰¹, E. D. Mendes Gouveia^{140a,140c}, L. Meng³⁶, X. T. Meng¹⁰⁶, S. Menke¹¹⁵, E. Meoni^{41a,41b}, S. Mergelmeyer¹⁹, S. A. M. Merkt¹³⁹, C. Merlassino²⁰, P. Mermod⁵⁴, L. Merola^{70a,70b}, C. Meroni^{69a}, G. Merz¹⁰⁶, O. Meshkov^{111,113}, J. K. R. Meshreki¹⁵¹, A. Messina^{73a,73b}, J. Metcalfe⁶, A. S. Mete¹⁷¹, C. Meyer⁶⁶, J. Meyer¹⁶⁰, J.-P. Meyer¹⁴⁵, H. Meyer Zu Theenhausen^{61a}, F. Miano¹⁵⁶, M. Michetti¹⁹, R. P. Middleton¹⁴⁴, L. Mijović⁵⁰, G. Mikenberg¹⁸⁰, M. Mikestikova¹⁴¹, M. Mikuz⁹², H. Mildner¹⁴⁹, M. Milesi¹⁰⁵, A. Milic¹⁶⁷, D. A. Millar⁹³, D. W. Miller³⁷, A. Milov¹⁸⁰, D. A. Milstead^{45a,45b}, R. A. Mina¹⁵³, A. A. Minaenko¹²³, M. Miñano Moya¹⁷⁴, I. A. Minashvili^{159b}, A. I. Mincer¹²⁵, B. Mindur^{84a}, M. Mineev⁸⁰, Y. Minegishi¹⁶³, L. M. Mir¹⁴, A. Mirto^{68a,68b}, K. P. Mistry¹³⁷, T. Mitani¹⁷⁹, J. Mitrevski¹¹⁴, V. A. Mitsou¹⁷⁴, M. Mittal^{60c}, O. Miu¹⁶⁷, A. Miucci²⁰, P. S. Miyagawa¹⁴⁹, A. Mizukami⁸², J. U. Mjörnmark⁹⁷, T. Mkrtychyan¹⁸⁴, M. Mlynarikova¹⁴³, T. Moa^{45a,45b}, K. Mochizuki¹¹⁰, P. Mogg⁵², S. Mohapatra³⁹, R. Moles-Valls²⁴, M. C. Mondragon¹⁰⁷, K. Mönig⁴⁶, J. Monk⁴⁰, E. Monnier¹⁰², A. Montalbano¹⁵², J. Montejo Berlingen³⁶, M. Montella⁹⁵, F. Monticelli⁸⁹, S. Monzani^{69a}, N. Morange⁶⁵, D. Moreno^{22a}, M. Moreno Llácer¹⁷⁴, C. Moreno Martinez¹⁴, P. Moretti^{55b}, M. Morgenstern¹²⁰, S. Morgenstern⁴⁸, D. Mori¹⁵², M. Morii⁵⁹, M. Morinaga¹⁷⁹, V. Morisbak¹³⁴, A. K. Morley³⁶, G. Mornacchi³⁶, A. P. Morris⁹⁵, L. Morvaj¹⁵⁵, P. Moschovakos³⁶, B. Moser¹²⁰, M. Mosidze^{159b}, T. Moskalets¹⁴⁵, H. J. Moss¹⁴⁹, J. Moss^{31.m}, E. J. W. Moyses¹⁰³, S. Muanza¹⁰², J. Mueller¹³⁹, R. S. P. Mueller¹¹⁴, D. Muenstermann⁹⁰, G. A. Mullier⁹⁷, D. P. Mungo^{69a,69b}, J. L. Munoz Martinez¹⁴, F. J. Munoz Sanchez¹⁰¹, P. Murin^{28b}, W. J. Murray^{178,144}, A. Murrone^{69a,69b}, M. Muškinja¹⁸, C. Mwewa^{33a}, A. G. Myagkov^{123.an}, A. A. Myers¹³⁹, J. Myers¹³², M. Myska¹⁴², B. P. Nachman¹⁸, O. Nackenhorst⁴⁷, A. Nag Nag⁴⁸, K. Nagai¹³⁵, K. Nagano⁸², Y. Nagasaka⁶², M. Nagel⁵², J. L. Nagle²⁹, E. Nagy¹⁰², A. M. Nairz³⁶, Y. Nakahama¹¹⁷, K. Nakamura⁸², T. Nakamura¹⁶³, I. Nakano¹²⁸, H. Nanjo¹³³, F. Napolitano^{61a}, R. F. Naranjo Garcia⁴⁶, R. Narayan⁴², I. Naryshkin¹³⁸, T. Naumann⁴⁶, G. Navarro^{22a}, P. Y. Nechaeva¹¹¹, F. Nechansky⁴⁶, T. J. Neep²¹,

A. Negri^{71a,71b}, M. Negrini^{23b}, C. Nellist⁵³, M. E. Nelson¹³⁵, S. Nemecek¹⁴¹, P. Nemethy¹²⁵, M. Nessi^{36,d}, M. S. Neubauer¹⁷³, M. Neumann¹⁸², R. Newhouse¹⁷⁵, P. R. Newman²¹, Y. S. Ng¹⁹, Y. W. Y. Ng¹⁷¹, B. Ngair^{35c}, H. D. N. Nguyen¹⁰², T. Nguyen Manh¹¹⁰, E. Nibigira³⁸, R. B. Nickerson¹³⁵, R. Nicolaidou¹⁴⁵, D. S. Nielsen⁴⁰, J. Nielsen¹⁴⁶, N. Nikiforou¹¹, V. Nikolaenko^{123,an}, I. Nikolic-Audit¹³⁶, K. Nikolopoulos²¹, P. Nilsson²⁹, H. R. Nindhito⁵⁴, Y. Ninomiya⁸², A. Nisati^{73a}, N. Nishu^{60c}, R. Nisius¹¹⁵, I. Nitsche⁴⁷, T. Nitta¹⁷⁹, T. Nobe¹⁶³, Y. Noguchi⁸⁶, I. Nomidis¹³⁶, M. A. Nomura²⁹, M. Nordberg³⁶, N. Norjoharuddeen¹³⁵, T. Novak⁹², O. Novgorodova⁴⁸, R. Novotny¹⁴², L. Nozka¹³¹, K. Ntekas¹⁷¹, E. Nurse⁹⁵, F. G. Oakham^{34,au}, H. Oberlack¹¹⁵, J. Ocariz¹³⁶, A. Ochi⁸³, I. Ochoa³⁹, J. P. Ochoa-Ricoux^{147a}, K. O'Connor²⁶, S. Oda⁸⁸, S. Odaka⁸², S. Oerdek⁵³, A. Ogrodnik^{84a}, A. Oh¹⁰¹, S. H. Oh⁴⁹, C. C. Ohm¹⁵⁴, H. Oide¹⁶⁵, M. L. Ojeda¹⁶⁷, H. Okawa¹⁶⁹, Y. Okazaki⁸⁶, M. W. O'Keefe⁹¹, Y. Okumura¹⁶³, T. Okuyama⁸², A. Olariu^{27b}, L. F. Oleiro Seabra^{140a}, S. A. Olivares Pino^{147a}, D. Oliveira Damazio²⁹, J. L. Oliver¹, M. J. R. Olsson¹⁷¹, A. Olszewski⁸⁵, J. Olszowska⁸⁵, D. C. O'Neil¹⁵², A. P. O'Neill¹³⁵, A. Onofre^{140a,140e}, P. U. E. Onyisi¹¹, H. Oppen¹³⁴, M. J. Oreglia³⁷, G. E. Orellana⁸⁹, D. Orestano^{75a,75b}, N. Orlando¹⁴, R. S. Orr¹⁶⁷, V. O'Shea⁵⁷, R. Ospanov^{60a}, G. Otero y Garzon³⁰, H. Otono⁸⁸, P. S. Ott^{61a}, M. Ouchrif^{35d}, J. Ouellette²⁹, F. Ould-Saada¹³⁴, A. Ouraou¹⁴⁵, Q. Ouyang^{15a}, M. Owen⁵⁷, R. E. Owen²¹, V. E. Ozcan^{12c}, N. Ozturk⁸, J. Pacalt¹³¹, H. A. Pacey³², K. Pachal⁴⁹, A. Pacheco Pages¹⁴, C. Padilla Aranda¹⁴, S. Pagan Griso¹⁸, M. Paganini¹⁸³, G. Palacino⁶⁶, S. Palazzo⁵⁰, S. Palestini³⁶, M. Palka^{84b}, D. Pallin³⁸, I. Panagoulas¹⁰, C. E. Pandini³⁶, J. G. Panduro Vazquez⁹⁴, P. Pani⁴⁶, G. Panizzo^{67a,67c}, L. Paolozzi⁵⁴, C. Papadatos¹¹⁰, K. Papageorgiou^{9,g}, S. Parajuli⁴³, A. Paramonov⁶, D. Paredes Hernandez^{63b}, S. R. Paredes Saenz¹³⁵, B. Parida¹⁶⁶, T. H. Park¹⁶⁷, A. J. Parker³¹, M. A. Parker³², F. Parodi^{55a,55b}, E. W. Parrish¹²¹, J. A. Parsons³⁹, U. Parzefall⁵², L. Pascual Dominguez¹³⁶, V. R. Pascuzzi¹⁶⁷, J. M. P. Pasner¹⁴⁶, F. Pasquali¹²⁰, E. Pasqualucci^{73a}, S. Passaggio^{55b}, F. Pastore⁹⁴, P. Pasuwan^{45a,45b}, S. Pataria¹⁰⁰, J. R. Pater¹⁰¹, A. Pathak^{181,i}, T. Pauly³⁶, B. Pearson¹¹⁵, M. Pedersen¹³⁴, L. Pedraza Diaz¹¹⁹, R. Pedro^{140a}, T. Peiffer⁵³, S. V. Peleganchuk^{122b,122a}, O. Penc¹⁴¹, H. Peng^{60a}, B. S. Peralva^{81a}, M. M. Perego⁶⁵, A. P. Pereira Peixoto^{140a}, D. V. Perepelitsa²⁹, F. Peri¹⁹, L. Perini^{69a,69b}, H. Pernegger³⁶, S. Perrella^{70a,70b}, K. Peters⁴⁶, R. F. Y. Peters¹⁰¹, B. A. Petersen³⁶, T. C. Petersen⁴⁰, E. Petit¹⁰², A. Petridis¹, C. Petridou¹⁶², P. Petroff⁶⁵, M. Petrov¹³⁵, F. Petrucci^{75a,75b}, M. Pettee¹⁸³, N. E. Pettersson¹⁰³, K. Petukhova¹⁴³, A. Peyaud¹⁴⁵, R. Pezoa^{147d}, L. Pezzotti^{71a,71b}, T. Pham¹⁰⁵, F. H. Phillips¹⁰⁷, P. W. Phillips¹⁴⁴, M. W. Phipps¹⁷³, G. Piacquadio¹⁵⁵, E. Pianori¹⁸, A. Picazio¹⁰³, R. H. Pickles¹⁰¹, R. Piegaia³⁰, D. Pietreanu^{27b}, J. E. Pilcher³⁷, A. D. Pilkington¹⁰¹, M. Pinamonti^{74a,74b}, J. L. Pinfold³, M. Pitt¹⁶¹, L. Pizzimento^{74a,74b}, M.-A. Pleier²⁹, V. Pleskot¹⁴³, E. Plotnikova⁸⁰, P. Podberezko^{122a,122b}, R. Poettgen⁹⁷, R. Poggi⁵⁴, L. Poggioli⁶⁵, I. Pogrebnayak¹⁰⁷, D. Pohl²⁴, I. Pokharel⁵³, G. Polesello^{71a}, A. Poley¹⁸, A. Policicchio^{73a,73b}, R. Polifka¹⁴³, A. Polini^{23b}, C. S. Pollard⁴⁶, V. Polychronakos²⁹, D. Ponomarenko¹¹², L. Pontecorvo³⁶, S. Popa^{27a}, G. A. Popeneciu^{27d}, L. Portales⁵, D. M. Portillo Quintero⁵⁸, S. Pospisil¹⁴², K. Potamianos⁴⁶, I. N. Potrap⁸⁰, C. J. Potter³², H. Potti¹¹, T. Poulsen⁹⁷, J. Poveda³⁶, T. D. Powell¹⁴⁹, G. Pownall⁴⁶, M. E. Pozo Astigarraga³⁶, P. Pralavorio¹⁰², S. Prell⁷⁹, D. Price¹⁰¹, M. Primavera^{68a}, S. Prince¹⁰⁴, M. L. Proffitt¹⁴⁸, N. Proklova¹¹², K. Prokofiev^{63c}, F. Prokoshin⁸⁰, S. Protopopescu²⁹, J. Proudfoot⁶, M. Przybycien^{84a}, D. Pudza¹³⁸, A. Puri¹⁷³, P. Puzo⁶⁵, J. Qian¹⁰⁶, Y. Qin¹⁰¹, A. Quadt⁵³, M. Queitsch-Maitland³⁶, A. Qureshi¹, M. Racko^{28a}, P. Rados¹⁰⁵, F. Ragusa^{69a,69b}, G. Rahal⁹⁸, J. A. Raine⁵⁴, S. Rajagopalan²⁹, A. Ramirez Morales⁹³, K. Ran^{15a,15d}, T. Rashid⁶⁵, S. Raspopov⁵, D. M. Rauch⁴⁶, F. Rauscher¹¹⁴, S. Rave¹⁰⁰, B. Ravina¹⁴⁹, I. Ravinovich¹⁸⁰, J. H. Rawling¹⁰¹, M. Raymond³⁶, A. L. Read¹³⁴, N. P. Readoff⁵⁸, M. Reale^{68a,68b}, D. M. Rebuffi^{71a,71b}, A. Redelbach¹⁷⁷, G. Redlinger²⁹, K. Reeves⁴³, L. Rehnisch¹⁹, J. Reichert¹³⁷, D. Reikher¹⁶¹, A. Reiss¹⁰⁰, A. Rej¹⁵¹, C. Rembser³⁶, M. Renda^{27b}, M. Rescigno^{73a}, S. Resconi^{69a}, E. D. Resseguie¹³⁷, S. Rettie¹⁷⁵, B. Reynolds¹²⁷, E. Reynolds²¹, O. L. Rezanova^{122a,122b}, P. Reznicek¹⁴³, E. Ricci^{76a,76b}, R. Richter¹¹⁵, S. Richter⁴⁶, E. Richter-Was^{84b}, O. Ricken²⁴, M. Ridet¹³⁶, P. Rieck¹¹⁵, O. Rifki⁴⁶, M. Rijssenbeek¹⁵⁵, A. Rimoldi^{71a,71b}, M. Rimoldi⁴⁶, L. Rinaldi^{23b}, G. Ripellino¹⁵⁴, I. Riu¹⁴, J. C. Rivera Vergara¹⁷⁶, F. Rizatdinova¹³⁰, E. Rizvi⁹³, C. Rizzi³⁶, R. T. Roberts¹⁰¹, S. H. Robertson^{104,ae}, M. Robin⁴⁶, D. Robinson³², J. E. M. Robinson⁴⁶, C. M. Robles Gajardo^{147d}, A. Robson⁵⁷, A. Rocchi^{74a,74b}, E. Rocco¹⁰⁰, C. Roda^{72a,72b}, S. Rodriguez Bosca¹⁷⁴, A. Rodriguez Perez¹⁴, D. Rodriguez Rodriguez¹⁷⁴, A. M. Rodriguez Vera^{168b}, S. Roe³⁶, O. Röhne¹³⁴, R. Röhrig¹¹⁵, R. A. Rojas^{147d}, C. P. A. Roland⁶⁶, J. Roloff²⁹, A. Romaniouk¹¹², M. Romano^{23a,23b}, N. Rompotis⁹¹, M. Ronzani¹²⁵, L. Roos¹³⁶, S. Rosati^{73a}, K. Rosbach⁵², G. Rosin¹⁰³, B. J. Rosser¹³⁷, E. Rossi⁴⁶, E. Rossi^{75a,75b}, E. Rossi^{70a,70b}, L. P. Rossi^{55b}, L. Rossini^{69a,69b}, R. Rosten¹⁴, M. Rotaru^{27b}, J. Rothberg¹⁴⁸, D. Rousseau⁶⁵, G. Rovelli^{71a,71b}

A. Roy¹¹, D. Roy^{33e}, A. Rozanov¹⁰², Y. Rozen¹⁶⁰, X. Ruan^{33e}, F. Rühr⁵², A. Ruiz-Martinez¹⁷⁴,
 A. Rummeler³⁶, Z. Rurikova⁵², N. A. Rusakovich⁸⁰, H. L. Russell¹⁰⁴, L. Rustige^{38,47}, J. P. Rutherford⁷,
 E. M. Rüttinger¹⁴⁹, M. Rybar³⁹, G. Rybkin⁶⁵, E. B. Rye¹³⁴, A. Ryzhov¹²³, J. A. Sabater Iglesias⁴⁶,
 P. Sabatini⁵³, G. Sabato¹²⁰, S. Sacerdoti⁶⁵, H.F.-W. Sadrozinski¹⁴⁶, R. Sadykov⁸⁰, F. Safai Tehrani^{73a},
 B. Safarzadeh Samani¹⁵⁶, P. Saha¹²¹, S. Saha¹⁰⁴, M. Sahinsoy^{61a}, A. Sahu¹⁸², M. Saimpert⁴⁶, M. Saito¹⁶³,
 T. Saito¹⁶³, H. Sakamoto¹⁶³, A. Sakharov^{125,am}, D. Salamani⁵⁴, G. Salamanna^{75a,75b}, J. E. Salazar Loyola^{147d},
 A. Salnikov¹⁵³, J. Salt¹⁷⁴, D. Salvatore^{41a,41b}, F. Salvatore¹⁵⁶, A. Salvucci^{63a,63b,63c}, A. Salzburger³⁶,
 J. Samarati³⁶, D. Sammel⁵², D. Sampsonidis¹⁶², D. Sampsonidou¹⁶², J. Sánchez¹⁷⁴, A. Sanchez Pineda^{36,67a,67c},
 H. Sandaker¹³⁴, C. O. Sander⁴⁶, I. G. Sanderswood⁹⁰, M. Sandhoff¹⁸², C. Sandoval^{22a}, D. P. C. Sankey¹⁴⁴,
 M. Sannino^{55a,55b}, Y. Sano¹¹⁷, A. Sansoni⁵¹, C. Santoni³⁸, H. Santos^{140a,140b}, S. N. Santpur¹⁸,
 A. Santra¹⁷⁴, A. Saponov⁸⁰, J. G. Saraiva^{140a,140d}, O. Sasaki⁸², K. Sato¹⁶⁹, F. Sauerburger⁵², E. Sauvan⁵,
 P. Savard^{167,au}, N. Savic¹¹⁵, R. Sawada¹⁶³, C. Sawyer¹⁴⁴, L. Sawyer^{96,ak}, C. Sbarra^{23b}, A. Sbrizzi^{23a},
 T. Scanlon⁹⁵, J. Schaarschmidt¹⁴⁸, P. Schacht¹¹⁵, B. M. Schachtner¹¹⁴, D. Schaefer³⁷, L. Schaefer¹³⁷,
 J. Schaeffer¹⁰⁰, S. Schaepe³⁶, U. Schäfer¹⁰⁰, A. C. Schaffer⁶⁵, D. Schaile¹¹⁴, R. D. Schamberger¹⁵⁵,
 E. Schanet¹¹⁴, N. Scharmberg¹⁰¹, V. A. Schegelsky¹³⁸, D. Scheirich¹⁴³, F. Schenck¹⁹, M. Schernau¹⁷¹,
 C. Schiavi^{55a,55b}, S. Schier¹⁴⁶, L. K. Schildgen²⁴, Z. M. Schillaci²⁶, E. J. Schioppa³⁶, M. Schioppa^{41b,41a},
 K. E. Schleicher⁵², S. Schlenker³⁶, K. R. Schmidt-Sommerfeld¹¹⁵, K. Schmieden³⁶, C. Schmitt¹⁰⁰,
 S. Schmitt⁴⁶, S. Schmitz¹⁰⁰, J. C. Schmoedel⁴⁶, U. Schnoor⁵², L. Schoeffel¹⁴⁵, A. Schoening^{61b},
 P. G. Scholer⁵², E. Schopf¹³⁵, M. Schott¹⁰⁰, J. F. P. Schouwenberg¹¹⁹, J. Schovancova³⁶, S. Schramm⁵⁴,
 F. Schroeder¹⁸², A. Schulte¹⁰⁰, H.-C. Schultz-Coulon^{61a}, M. Schumacher⁵², B. A. Schumm¹⁴⁶, Ph. Schune¹⁴⁵,
 A. Schwartzman¹⁵³, T. A. Schwarz¹⁰⁶, Ph. Schwemling¹⁴⁵, R. Schwienhorst¹⁰⁷, A. Sciandra¹⁴⁶, G. Sciolla²⁶,
 M. Scodiggio⁴⁶, M. Scornajenghi^{41a,41b}, F. Scuri^{72a}, F. Scutti¹⁰⁵, L. M. Scyboz¹¹⁵, C. D. Sebastiani^{73a,73b},
 P. Seema¹⁹, S. C. Seidel¹¹⁸, A. Seiden¹⁴⁶, B. D. Seidlitz²⁹, T. Seiss³⁷, J. M. Seixas^{81b}, G. Sekhniaidze^{70a},
 K. Sekhon¹⁰⁶, S. J. Sekula⁴², N. Semprini-Cesari^{23b,23a}, S. Sen⁴⁹, C. Serfon⁷⁷, L. Serin⁶⁵, L. Serkin^{67a,67b},
 M. Sessa^{60a}, H. Severini¹²⁹, T. Šfiligoj⁹², F. Sforza^{55a,55b}, A. Sfyrła⁵⁴, E. Shabalina⁵³, J. D. Shahinian¹⁴⁶,
 N. W. Shaikh^{45a,45b}, D. Shaked Renous¹⁸⁰, L. Y. Shan^{15a}, J. T. Shank²⁵, M. Shapiro¹⁸, A. Sharma¹³⁵,
 A. S. Sharma¹, P. B. Shatalov¹²⁴, K. Shaw¹⁵⁶, S. M. Shaw¹⁰¹, A. Shcherbakova¹³⁸, M. Shehade¹⁸⁰,
 Y. Shen¹²⁹, A. D. Sherman²⁵, P. Sherwood⁹⁵, L. Shi^{158,as}, S. Shimizu⁸², C. O. Shimmin¹⁸³, Y. Shimogama¹⁷⁹,
 M. Shimojima¹¹⁶, I. P. J. Shipsey¹³⁵, S. Shirabe⁸⁸, M. Shiyakova^{80,ac}, J. Shlomi¹⁸⁰, A. Shmeleva¹¹¹,
 M. J. Shochet³⁷, J. Shojaii¹⁰⁵, D. R. Shope¹²⁹, S. Shrestha¹²⁷, E. M. Shrif^{33e}, E. Shulga¹⁸⁰, P. Sicho¹⁴¹,
 A. M. Sickles¹⁷³, P. E. Sidebo¹⁵⁴, E. Sideras Haddad^{33e}, O. Sidiropoulou³⁶, A. Sidoti^{23a,23b}, F. Siegert⁴⁸,
 Dj. Sijacki¹⁶, M. Jr. Silva¹⁸¹, M. V. Silva Oliveira^{81a}, S. B. Silverstein^{45a}, S. Simion⁶⁵, E. Simioni¹⁰⁰,
 R. Simoniello¹⁰⁰, S. Simsek^{12b}, P. Sinervo¹⁶⁷, V. Sinetckii^{113,111}, N. B. Sinev¹³², M. Sioli^{23a,23b}, I. Siral¹³²,
 S. Yu. Sivoklov¹¹³, J. Sjölin^{45a,45b}, E. Skorda⁹⁷, P. Skubic¹²⁹, M. Slawinska⁸⁵, K. Sliwa¹⁷⁰, R. Slovak¹⁴³,
 V. Smakhtin¹⁸⁰, B. H. Smart¹⁴⁴, J. Smiesko^{28a}, N. Smirnov¹¹², S. Yu. Smirnov¹¹², Y. Smirnov¹¹²,
 L. N. Smirnova^{113,u}, O. Smirnova⁹⁷, J. W. Smith⁵³, M. Smizanska⁹⁰, K. Smolek¹⁴², A. Smykiewicz⁸⁵,
 A. A. Snesarev¹¹¹, H. L. Snoek¹²⁰, I. M. Snyder¹³², S. Snyder²⁹, R. Sobie^{176,ae}, A. Soffer¹⁶¹, A. Sjøgaard⁵⁰,
 F. Sohns⁵³, C. A. Solans Sanchez³⁶, E. Yu. Soldatov¹¹², U. Soldevila¹⁷⁴, A. A. Solodkov¹²³, A. Soloshenko⁸⁰,
 O. V. Solovyov¹²³, V. Solov'yev¹³⁸, P. Sommer¹⁴⁹, H. Son¹⁷⁰, W. Song¹⁴⁴, W. Y. Song^{168b},
 A. Sopczak¹⁴², F. Sopkova^{28b}, C. L. Sotiropoulou^{72a,72b}, S. Sottocornola^{71a,71b}, R. Soualah^{67a,67c,f},
 A. M. Soukharev^{122b,122a}, D. South⁴⁶, S. Spagnolo^{68a,68b}, M. Spalla¹¹⁵, M. Spangenberg¹⁷⁸, F. Spanò⁹⁴,
 D. Sperlich⁵², T. M. Spieker^{61a}, R. Spighi^{23b}, G. Spigo³⁶, M. Spina¹⁵⁶, D. P. Spiteri⁵⁷, M. Spousta¹⁴³,
 A. Stabile^{69a,69b}, B. L. Stamas¹²¹, R. Stamen^{61a}, M. Stamenkovic¹²⁰, E. Stanecka⁸⁵, B. Stanislaus¹³⁵,
 M. M. Stanitzki⁴⁶, M. Stankaityte¹³⁵, B. Stapf¹²⁰, E. A. Starchenko¹²³, G. H. Stark¹⁴⁶, J. Stark⁵⁸,
 S. H. Stark⁴⁰, P. Staroba¹⁴¹, P. Starovoitov^{61a}, S. Stärz¹⁰⁴, R. Staszewski⁸⁵, G. Stavropoulos⁴⁴,
 M. Stegler⁴⁶, P. Steinberg²⁹, A. L. Steinhebel¹³², B. Stelzer¹⁵², H. J. Stelzer¹³⁹, O. Stelzer-Chilton^{168a},
 H. Stenzel⁵⁶, T. J. Stevenson¹⁵⁶, G. A. Stewart³⁶, M. C. Stockton³⁶, G. Stoicea^{27b}, M. Stolarski^{140a},
 S. Stonjek¹¹⁵, A. Straessner⁴⁸, J. Strandberg¹⁵⁴, S. Strandberg^{45a,45b}, M. Strauss¹²⁹, P. Striznec^{28b},
 R. Ströhmer¹⁷⁷, D. M. Strom¹³², R. Stroynowski⁴², A. Strubig⁵⁰, S. A. Stucci²⁹, B. Stugu¹⁷,
 J. Stupak¹²⁹, N. A. Styles⁴⁶, D. Su¹⁵³, S. Suchek^{61a}, V. V. Sulin¹¹¹, M. J. Sullivan⁹¹, D. M. S. Sultan⁵⁴,
 S. Sultansoy^{4c}, T. Sumida⁸⁶, S. Sun¹⁰⁶, X. Sun³, K. Suruliz¹⁵⁶, C. J. E. Suster¹⁵⁷, M. R. Sutton¹⁵⁶,
 S. Suzuki⁸², M. Svatos¹⁴¹, M. Swiatlowski³⁷, S. P. Swift², T. Swirski¹⁷⁷, A. Sydorenko¹⁰⁰, I. Sykora^{28a},
 M. Sykora¹⁴³, T. Sykora¹⁴³, D. Ta¹⁰⁰, K. Tackmann^{46,aa}, J. Taenzer¹⁶¹, A. Taffard¹⁷¹, R. Tafirout^{168a},

H. Takai²⁹, R. Takashima⁸⁷, K. Takeda⁸³, T. Takeshita¹⁵⁰, E. P. Takeva⁵⁰, Y. Takubo⁸², M. Talby¹⁰², A. A. Talyshev^{122b,122a}, N. M. Tamir¹⁶¹, J. Tanaka¹⁶³, M. Tanaka¹⁶⁵, R. Tanaka⁶⁵, S. Tapia Araya¹⁷³, S. Tapprogge¹⁰⁰, A. Tarek Abouelfadl Mohamed¹³⁶, S. Tarem¹⁶⁰, K. Tariq^{60b}, G. Tarna^{27b,c}, G. F. Tartarelli^{69a}, P. Tas¹⁴³, M. Tasevsky¹⁴¹, T. Tashiro⁸⁶, E. Tassi^{41b,41a}, A. Tavares Delgado^{140a,140b}, Y. Tayalati^{35c}, A. J. Taylor⁵⁰, G. N. Taylor¹⁰⁵, W. Taylor^{168b}, A. S. Tee⁹⁰, R. Teixeira De Lima¹⁵³, P. Teixeira-Dias⁹⁴, H. Ten Kate³⁶, J. J. Teoh¹²⁰, S. Terada⁸², K. Terashi¹⁶³, J. Terron⁹⁹, S. Terzo¹⁴, M. Testa⁵¹, R. J. Teuscher^{167,ac}, S. J. Thais¹⁸³, T. Theveneaux-Pelzer⁴⁶, F. Thiele⁴⁰, D. W. Thomas⁹⁴, J. O. Thomas⁴², J. P. Thomas²¹, A. S. Thompson⁵⁷, P. D. Thompson²¹, L. A. Thomsen¹⁸³, E. Thomson¹³⁷, E. J. Thorpe⁹³, R. E. Ticse Torres⁵³, V. O. Tikhomirov^{111,ao}, Yu. A. Tikhonov^{122a,122b}, S. Timoshenko¹¹², P. Tipton¹⁸³, S. Tisserant¹⁰², K. Todome^{23a,23b}, S. Todorova-Nova⁵, S. Todt⁴⁸, J. Tojo⁸⁸, S. Tokár^{28a}, K. Tokushuku⁸², E. Tolley¹²⁷, K. G. Tomiwa^{33e}, M. Tomoto¹¹⁷, L. Tompkins^{153,p}, B. Tong⁵⁹, P. Tornambe¹⁰³, E. Torrence¹³², H. Torres⁴⁸, E. Torró Pastor¹⁴⁸, C. Toscirri¹³⁵, J. Toth^{102,ad}, D. R. Tovey¹⁴⁹, A. Traet¹⁷, C. J. Treado¹²⁵, T. Trefzger¹⁷⁷, F. Tresoldi¹⁵⁶, A. Tricoli²⁹, I. M. Trigger^{168a}, S. Trincaz-Duvoid¹³⁶, D. A. Trischuk¹⁷⁵, W. Trischuk¹⁶⁷, B. Trocmé⁵⁸, A. Trofymov¹⁴⁵, C. Troncon^{69a}, M. Trovatielli¹⁷⁶, F. Trovato¹⁵⁶, L. Truong^{33c}, M. Trzebinski⁸⁵, A. Trzupke⁸⁵, F. Tsai⁴⁶, J.C.-L. Tseng¹³⁵, P. V. Tsireshka^{108,aj}, A. Tsigotis^{162,x}, V. Tsiskaridze¹⁵⁵, E. G. Tskhadadze^{159a}, M. Tsopoulou¹⁶², I. I. Tsukerman¹²⁴, V. Tsulaia¹⁸, S. Tsuno⁸², D. Tsybychev¹⁵⁵, Y. Tu^{63b}, A. Tudorache^{27b}, V. Tudorache^{27b}, T. T. Tulbure^{27a}, A. N. Tuna⁵⁹, S. Turchikhin⁸⁰, D. Turgeman¹⁸⁰, I. Turk Cakir^{4b,v}, R. J. Turner²¹, R. T. Turra^{69a}, P. M. Tuts³⁹, S. Tzamaras¹⁶², E. Tzovara¹⁰⁰, G. Uccielli⁴⁷, K. Uchida¹⁶³, I. Ueda⁸², M. Ughetto^{45a,45b}, F. Ukegawa¹⁶⁹, G. Unal³⁶, A. Undrus²⁹, G. Unel¹⁷¹, F. C. Ungaro¹⁰⁵, Y. Unno⁸², K. Uno¹⁶³, J. Urban^{28b}, P. Urquijo¹⁰⁵, G. Usai⁸, Z. Uysal^{12d}, V. Vacek¹⁴², B. Vachon¹⁰⁴, K. O. H. Vadla¹³⁴, A. Vaidya⁹⁵, C. Valderanis¹¹⁴, E. Valdes Santurio^{45a,45b}, M. Valente⁵⁴, S. Valentini^{23a,23b}, A. Valero¹⁷⁴, L. Valéry⁴⁶, R. A. Vallance²¹, A. Vallier³⁶, J. A. Valls Ferrer¹⁷⁴, T. R. Van Daalen¹⁴, P. Van Gemmeren⁶, I. Van Vulpen¹²⁰, M. Vanadia^{74a,74b}, W. Vandelli³⁶, E. R. Vandewall¹³⁰, A. Vaniachine¹⁶⁶, D. Vannicola^{73a,73b}, R. Vari^{73a}, E. W. Varnes⁷, C. Varni^{55a,55b}, T. Varol¹⁵⁸, D. Varouchas⁶⁵, K. E. Varvell¹⁵⁷, M. E. Vasile^{27b}, G. A. Vasquez¹⁷⁶, J. G. Vasquez¹⁸³, F. Vazeille³⁸, D. Vazquez Furelos¹⁴, T. Vazquez Schroeder³⁶, J. Veatch⁵³, V. Vecchio^{75a,75b}, M. J. Veen¹²⁰, L. M. Veloce¹⁶⁷, F. Veloso^{140a,140c}, S. Veneziano^{73a}, A. Ventura^{68a,68b}, N. Venturi³⁶, A. Verbytskyi¹¹⁵, V. Vercesi^{71a}, M. Verducci^{72a,72b}, C. M. Vergel Infante⁷⁹, C. Vergis²⁴, W. Verkerke¹²⁰, A. T. Vermeulen¹²⁰, J. C. Vermeulen¹²⁰, M. C. Vetterli^{152,au}, N. Viaux Maira^{147d}, M. Vicente Barreto Pinto⁵⁴, T. Vickey¹⁴⁹, O. E. Vickey Boeriu¹⁴⁹, G. H. A. Viehhauser¹³⁵, L. Vigani^{61b}, M. Villa^{23b,23a}, M. Villaplana Perez^{69a,69b}, E. Vilucchi⁵¹, M. G. Vincker³⁴, G. S. Virdee²¹, A. Vishwakarma⁴⁶, C. Vittori^{23a,23b}, I. Vivarelli¹⁵⁶, M. Vogel¹⁸², P. Vokac¹⁴², S. E. von Buddenbrock^{33e}, E. Von Toerne²⁴, V. Vorobel¹⁴³, K. Vorobev¹¹², M. Vos¹⁷⁴, J. H. Vosseveld⁹¹, M. Vozak¹⁰¹, N. Vranjes¹⁶, M. Vranjes Milosavljevic¹⁶, V. Vrba¹⁴², M. Vreeswijk¹²⁰, R. Vuillermet³⁶, I. Vukotic³⁷, P. Wagner²⁴, W. Wagner¹⁸², J. Wagner-Kuhr¹¹⁴, S. Wahdan¹⁸², H. Wahlberg⁸⁹, V. M. Walbrecht¹¹⁵, J. Walder⁹⁰, R. Walker¹¹⁴, S. D. Walker⁹⁴, W. Walkowiak¹⁵¹, V. Wallangen^{45a,45b}, A. M. Wang⁵⁹, C. Wang^{60c}, C. Wang^{60b}, F. Wang¹⁸¹, H. Wang¹⁸, H. Wang³, J. Wang^{63a}, J. Wang¹⁵⁷, J. Wang^{61b}, P. Wang⁴², Q. Wang¹²⁹, R.-J. Wang¹⁰⁰, R. Wang^{60a}, R. Wang⁶, S. M. Wang¹⁵⁸, W. T. Wang^{60a}, W. Wang^{15c}, W. X. Wang^{60a,af}, Y. Wang^{60a,al}, Z. Wang^{60c}, C. Wanotayaroj⁴⁶, A. Warburton¹⁰⁴, C. P. Ward³², D. R. Wardrope⁹⁵, N. Warrack⁵⁷, A. Washbrook⁵⁰, A. T. Watson²¹, M. F. Watson²¹, G. Watts¹⁴⁸, B. M. Waugh⁹⁵, A. F. Webb¹¹, S. Webb¹⁰⁰, C. Weber¹⁸³, M. S. Weber²⁰, S. A. Weber³⁴, S. M. Weber^{61a}, A. R. Weidberg¹³⁵, J. Weingarten⁴⁷, M. Weirich¹⁰⁰, C. Weiser⁵², P. S. Wells³⁶, T. Wenaus²⁹, T. Wengler³⁶, S. Wenig³⁶, N. Vermes²⁴, M. D. Werner⁷⁹, M. Wessels^{61a}, T. D. Weston²⁰, K. Whalen¹³², N. L. Whallon¹⁴⁸, A. M. Wharton⁹⁰, A. S. White¹⁰⁶, A. White⁸, M. J. White¹, D. Whiteson¹⁷¹, B. W. Whitmore⁹⁰, W. Wiedenmann¹⁸¹, M. Wieler¹⁴⁴, N. Wieseotte¹⁰⁰, C. Wiglesworth⁴⁰, L. A. M. Wiik-Fuchs⁵², F. Wilk¹⁰¹, H. G. Wilkens³⁶, L. J. Wilkins⁹⁴, H. H. Williams¹³⁷, S. Williams³², C. Willis¹⁰⁷, S. Willocq¹⁰³, J. A. Wilson²¹, I. Wingerter-Seetz⁵, E. Winkels¹⁵⁶, F. Winklmeier¹³², O. J. Winston¹⁵⁶, B. T. Winter⁵², M. Wittgen¹⁵³, M. Wobisch⁹⁶, A. Wolf¹⁰⁰, T. M. H. Wolf¹²⁰, R. Wolff¹⁰², R. Wölker¹³⁵, J. Wollrath⁵², M. W. Wolter⁸⁵, H. Wolters^{140a,140c}, V. W. S. Wong¹⁷⁵, N. L. Woods¹⁴⁶, S. D. Worm²¹, B. K. Wosiek⁸⁵, K. W. Woźniak⁸⁵, K. Wraight⁵⁷, S. L. Wu¹⁸¹, X. Wu⁵⁴, Y. Wu^{60a}, T. R. Wyatt¹⁰¹, B. M. Wynne⁵⁰, S. Xella⁴⁰, Z. Xi¹⁰⁶, L. Xia¹⁷⁸, X. Xiao¹⁰⁶, I. Xiotidis¹⁵⁶, D. Xu^{15a}, H. Xu^{60a,c}, L. Xu²⁹, T. Xu¹⁴⁵, W. Xu¹⁰⁶, Z. Xu^{60b}, Z. Xu¹⁵³, B. Yabsley¹⁵⁷, S. Yacoub^{33a}, K. Yajima¹³³, D. P. Yallup⁹⁵, D. Yamaguchi¹⁶⁵, Y. Yamaguchi¹⁶⁵, A. Yamamoto⁸², M. Yamatani¹⁶³, T. Yamazaki¹⁶³, Y. Yamazaki⁸³, Z. Yan²⁵, H. J. Yang^{60c,60d}, H. T. Yang¹⁸, S. Yang⁷⁸, X. Yang^{60b,58}, Y. Yang¹⁶³, W.-M. Yao¹⁸, Y. C. Yap⁴⁶, Y. Yasu⁸², E. Yatsenko^{60c,60d}, J. Ye⁴²

S. Ye²⁹ , I. Yeletsikh⁸⁰ , M. R. Yexley⁹⁰ , E. Yigitbasi²⁵ , K. Yorita¹⁷⁹ , K. Yoshihara¹³⁷ , C. J. S. Young³⁶ , C. Young¹⁵³ , J. Yu⁷⁹ , R. Yuan^{60b,h} , X. Yue^{61a} , S. P. Yuen²⁴ , M. Zaazoua^{35e} , B. Zabinski⁸⁵ , G. Zacharis¹⁰ , E. Zaffaroni⁵⁴ , J. Zahreddine¹³⁶ , A. M. Zaitsev^{123,an} , T. Zakareishvili^{159b} , N. Zakharchuk³⁴ , S. Zambito⁵⁹ , D. Zanzi³⁶ , D. R. Zaripovas⁵⁷ , S. V. Zeibner⁴⁷ , C. Zeitnitz¹⁸² , G. Zemaityte¹³⁵ , J. C. Zeng¹⁷³ , O. Zenin¹²³ , T. Ženiš^{28a} , D. Zerwas⁶⁵ , M. Zgubič¹³⁵ , B. Zhang^{15c} , D. F. Zhang^{15b} , G. Zhang^{15b} , H. Zhang^{15c} , J. Zhang⁶ , L. Zhang^{15c} , L. Zhang^{60a} , M. Zhang¹⁷³ , R. Zhang²⁴ , X. Zhang^{60b} , Y. Zhang^{15a,15d} , Z. Zhang^{63a} , Z. Zhang⁶⁵ , P. Zhao⁴⁹ , Y. Zhao^{60b} , Z. Zhao^{60a} , A. Zhemchugov⁸⁰ , Z. Zheng¹⁰⁶ , D. Zhong¹⁷³ , B. Zhou¹⁰⁶ , C. Zhou¹⁸¹ , M. S. Zhou^{15a,15d} , M. Zhou¹⁵⁵ , N. Zhou^{60c} , Y. Zhou⁷ , C. G. Zhu^{60b} , C. Zhu^{15a,15d} , H. L. Zhu^{60a} , H. Zhu^{15a} , J. Zhu¹⁰⁶ , Y. Zhu^{60a} , X. Zhuang^{15a} , K. Zhukov¹¹¹ , V. Zhulanov^{122a,122b} , D. Zieminska⁶⁶ , N. I. Zimine⁸⁰ , S. Zimmermann⁵² , Z. Zinonos¹¹⁵ , M. Ziolkowski¹⁵¹ , L. Živković¹⁶ , G. Zobernig¹⁸¹ , A. Zoccoli^{23a,23b} , K. Zoch⁵³ , T. G. Zorbas¹⁴⁹ , R. Zou³⁷ , L. Zwalinski³⁶ 

¹ Department of Physics, University of Adelaide, Adelaide, Australia

² Physics Department, SUNY Albany, Albany, NY, USA

³ Department of Physics, University of Alberta, Edmonton, AB, Canada

⁴ (a) Department of Physics, Ankara University, Ankara, Turkey; (b) Istanbul Aydin University, Istanbul, Turkey; (c) Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey

⁵ LAPP, Université Grenoble Alpes, Université Savoie Mont Blanc, CNRS/IN2P3, Annecy, France

⁶ High Energy Physics Division, Argonne National Laboratory, Argonne, IL, USA

⁷ Department of Physics, University of Arizona, Tucson, AZ, USA

⁸ Department of Physics, University of Texas at Arlington, Arlington, TX, USA

⁹ Physics Department, National and Kapodistrian University of Athens, Athens, Greece

¹⁰ Physics Department, National Technical University of Athens, Zografou, Greece

¹¹ Department of Physics, University of Texas at Austin, Austin, TX, USA

¹² (a) Bahcesehir University, Faculty of Engineering and Natural Sciences, Istanbul, Turkey; (b) Istanbul Bilgi University, Faculty of Engineering and Natural Sciences, Istanbul, Turkey; (c) Department of Physics, Bogazici University, Istanbul, Turkey; (d) Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey

¹³ Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan

¹⁴ Institut de Física d'Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona, Spain

¹⁵ (a) Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China; (b) Physics Department, Tsinghua University, Beijing, China; (c) Department of Physics, Nanjing University, Nanjing, China; (d) University of Chinese Academy of Science (UCAS), Beijing, China

¹⁶ Institute of Physics, University of Belgrade, Belgrade, Serbia

¹⁷ Department for Physics and Technology, University of Bergen, Bergen, Norway

¹⁸ Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, CA, USA

¹⁹ Institut für Physik, Humboldt Universität zu Berlin, Berlin, Germany

²⁰ Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland

²¹ School of Physics and Astronomy, University of Birmingham, Birmingham, UK

²² (a) Facultad de Ciencias y Centro de Investigaciones, Universidad Antonio Nariño, Bogotá, Colombia; (b) Departamento de Física, Universidad Nacional de Colombia, Bogotá, Colombia, Colombia

²³ (a) Dipartimento di Fisica, INFN Bologna and Università di Bologna, Bologna, Italy; (b) INFN Sezione di Bologna, Bologna, Italy

²⁴ Physikalisches Institut, Universität Bonn, Bonn, Germany

²⁵ Department of Physics, Boston University, Boston, MA, USA

²⁶ Department of Physics, Brandeis University, Waltham, MA, USA

²⁷ (a) Transilvania University of Brasov, Brasov, Romania; (b) Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest, Romania; (c) Department of Physics, Alexandru Ioan Cuza University of Iasi, Iasi, Romania; (d) Physics Department, National Institute for Research and Development of Isotopic and Molecular Technologies, Cluj-Napoca, Romania; (e) University Politehnica Bucharest, Bucharest, Romania; (f) West University in Timisoara, Timisoara, Romania

²⁸ (a) Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava, Slovakia; (b) Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic

- ²⁹ Physics Department, Brookhaven National Laboratory, Upton, NY, USA
- ³⁰ Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
- ³¹ California State University, CA, USA
- ³² Cavendish Laboratory, University of Cambridge, Cambridge, UK
- ³³ ^(a)Department of Physics, University of Cape Town, Cape Town, South Africa; ^(b)iThemba Labs, Western Cape, South Africa; ^(c)Department of Mechanical Engineering Science, University of Johannesburg, Johannesburg, South Africa ; ^(d)Department of Physics, University of South Africa, Pretoria, South Africa; ^(e)School of Physics, University of the Witwatersrand, Johannesburg, South Africa
- ³⁴ Department of Physics, Carleton University, Ottawa, ON, Canada
- ³⁵ ^(a)Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies - Université Hassan II, Casablanca, Morocco; ^(b)Faculté des Sciences, Université Ibn-Tofail, Kénitra, Morocco; ^(c)Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech, Marrakesh, Morocco; ^(d)Faculté des Sciences, Université Mohamed Premier and LTPPM, Oujda, Morocco; ^(e)Faculté des sciences, Université Mohammed V, Rabat, Morocco
- ³⁶ CERN, Geneva, Switzerland
- ³⁷ Enrico Fermi Institute, University of Chicago, Chicago, IL, USA
- ³⁸ LPC, Université Clermont Auvergne, CNRS/IN2P3, Clermont-Ferrand, France
- ³⁹ Nevis Laboratory, Columbia University, Irvington, NY, USA
- ⁴⁰ Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
- ⁴¹ ^(a)Dipartimento di Fisica, Università della Calabria, Rende, Italy; ^(b)INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati, Italy
- ⁴² Physics Department, Southern Methodist University, Dallas, TX, USA
- ⁴³ Physics Department, University of Texas at Dallas, Richardson, TX, USA
- ⁴⁴ National Centre for Scientific Research “Demokritos”, Agia Paraskevi, Greece
- ⁴⁵ ^(a)Department of Physics, Stockholm University, Stockholm, Sweden; ^(b)Oskar Klein Centre, Stockholm, Sweden
- ⁴⁶ Deutsches Elektronen-Synchrotron DESY, Hamburg and Zeuthen, Germany
- ⁴⁷ Lehrstuhl für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
- ⁴⁸ Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany
- ⁴⁹ Department of Physics, Duke University, Durham, NC, USA
- ⁵⁰ SUPA - School of Physics and Astronomy, University of Edinburgh, Edinburgh, UK
- ⁵¹ INFN e Laboratori Nazionali di Frascati, Frascati, Italy
- ⁵² Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg, Germany
- ⁵³ II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen, Germany
- ⁵⁴ Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève, Switzerland
- ⁵⁵ ^(a)Dipartimento di Fisica, Università di Genova, Genova, ; ^(b)INFN Sezione di Genova, Genoa, Italy
- ⁵⁶ II. Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany
- ⁵⁷ SUPA - School of Physics and Astronomy, University of Glasgow, Glasgow, UK
- ⁵⁸ LPSC, Université Grenoble Alpes, CNRS/IN2P3, Grenoble INP, Grenoble, France
- ⁵⁹ Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, MA, USA
- ⁶⁰ ^(a)Department of Modern Physics and State Key Laboratory of Particle Detection and Electronics, University of Science and Technology of China, Hefei, China; ^(b)Institute of Frontier and Interdisciplinary Science and Key Laboratory of Particle Physics and Particle Irradiation (MOE), Shandong University, Qingdao, China; ^(c)School of Physics and Astronomy, Shanghai Jiao Tong University, KLPPAC-MoE, SKLPPC, Shanghai, China; ^(d)Tsung-Dao Lee Institute, Shanghai, China
- ⁶¹ ^(a)Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany; ^(b)Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
- ⁶² Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan
- ⁶³ ^(a)Department of Physics, Chinese University of Hong Kong, Shatin, N.T., Hong Kong, China; ^(b)Department of Physics, University of Hong Kong, Hong Kong, China; ^(c)Department of Physics and Institute for Advanced Study, Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China
- ⁶⁴ Department of Physics, National Tsing Hua University, Hsinchu, Taiwan
- ⁶⁵ IJCLab, Université Paris-Saclay, CNRS/IN2P3, 91405, Orsay, France
- ⁶⁶ Department of Physics, Indiana University, Bloomington, IN, USA

- 67 (a) INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine, Italy; (b) ICTP, Trieste, Italy; (c) Dipartimento Politecnico di Ingegneria e Architettura, Università di Udine, Udine, Italy
- 68 (a) INFN Sezione di Lecce, Zona Monte, Italy; (b) Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy
- 69 (a) INFN Sezione di Milano, Milano, Italy; (b) Dipartimento di Fisica, Università di Milano, Milano, Italy
- 70 (a) INFN Sezione di Napoli, Naples, Italy; (b) Dipartimento di Fisica, Università di Napoli, Napoli, Italy
- 71 (a) INFN Sezione di Pavia, Pavia, Italy; (b) Dipartimento di Fisica, Università di Pavia, Pavia, Italy
- 72 (a) INFN Sezione di Pisa, Pisa, Italy; (b) Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy
- 73 (a) INFN Sezione di Roma, Rome, Italy; (b) Dipartimento di Fisica, Sapienza Università di Roma, Rome, Italy
- 74 (a) INFN Sezione di Roma Tor Vergata, Rome, Italy; (b) Dipartimento di Fisica, Università di Roma Tor Vergata, Rome, Italy
- 75 (a) INFN Sezione di Roma Tre, Rome, Italy; (b) Dipartimento di Matematica e Fisica, Università Roma Tre, Rome, Italy
- 76 (a) INFN-TIFPA, Povo, Italy; (b) Università degli Studi di Trento, Trento, Italy
- 77 Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
- 78 University of Iowa, Iowa City, IA, USA
- 79 Department of Physics and Astronomy, Iowa State University, Ames, IA, USA
- 80 Joint Institute for Nuclear Research, Dubna, Russia
- 81 (a) Departamento de Engenharia Elétrica, Universidade Federal de Juiz de Fora (UFJF), Juiz de Fora, Brazil; (b) Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro, Brazil; (c) Universidade Federal de São João del Rei (UFSJ), São João del Rei, Brazil; (d) Instituto de Física, Universidade de São Paulo, São Paulo, Brazil
- 82 KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
- 83 Graduate School of Science, Kobe University, Kobe, Japan
- 84 (a) AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland; (b) Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland
- 85 Institute of Nuclear Physics Polish Academy of Sciences, Krakow, Poland
- 86 Faculty of Science, Kyoto University, Kyoto, Japan
- 87 Kyoto University of Education, Kyoto, Japan
- 88 Research Center for Advanced Particle Physics and Department of Physics, Kyushu University, Fukuoka, Japan
- 89 Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina
- 90 Physics Department, Lancaster University, Lancaster, UK
- 91 Oliver Lodge Laboratory, University of Liverpool, Liverpool, UK
- 92 Department of Experimental Particle Physics, Jožef Stefan Institute and Department of Physics, University of Ljubljana, Ljubljana, Slovenia
- 93 School of Physics and Astronomy, Queen Mary University of London, London, UK
- 94 Department of Physics, Royal Holloway University of London, Egham, UK
- 95 Department of Physics and Astronomy, University College London, London, UK
- 96 Louisiana Tech University, Ruston, LA, USA
- 97 Fysiska institutionen, Lunds universitet, Lund, Sweden
- 98 Centre de Calcul de l'Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France
- 99 Departamento de Física Teórica C-15 and CIAFF, Universidad Autónoma de Madrid, Madrid, Spain
- 100 Institut für Physik, Universität Mainz, Mainz, Germany
- 101 School of Physics and Astronomy, University of Manchester, Manchester, UK
- 102 CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille, France
- 103 Department of Physics, University of Massachusetts, Amherst, MA, USA
- 104 Department of Physics, McGill University, Montreal, QC, Canada
- 105 School of Physics, University of Melbourne, Victoria, Australia
- 106 Department of Physics, University of Michigan, Ann Arbor, MI, USA
- 107 Department of Physics and Astronomy, Michigan State University, East Lansing, MI, USA
- 108 B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Belarus
- 109 Research Institute for Nuclear Problems of Byelorussian State University, Minsk, Belarus
- 110 Group of Particle Physics, University of Montreal, Montreal, QC, Canada
- 111 P.N. Lebedev Physical Institute of the Russian Academy of Sciences, Moscow, Russia

- 112 National Research Nuclear University MEPhI, Moscow, Russia
- 113 D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia
- 114 Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany
- 115 Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany
- 116 Nagasaki Institute of Applied Science, Nagasaki, Japan
- 117 Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan
- 118 Department of Physics and Astronomy, University of New Mexico, Albuquerque, NM, USA
- 119 Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, The Netherlands
- 120 Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, The Netherlands
- 121 Department of Physics, Northern Illinois University, DeKalb, IL, USA
- 122 (a) Budker Institute of Nuclear Physics and NSU, SB RAS, Novosibirsk, Russia; (b) Novosibirsk State University Novosibirsk, Novosibirsk, Russia
- 123 Institute for High Energy Physics of the National Research Centre Kurchatov Institute, Protvino, Russia
- 124 Institute for Theoretical and Experimental Physics named by A.I. Alikhanov of National Research Centre “Kurchatov Institute”, Moscow, Russia
- 125 Department of Physics, New York University, New York, NY, USA
- 126 Ochanomizu University, Otsuka, Bunkyo-ku, Tokyo, Japan
- 127 Ohio State University, Columbus, OH, USA
- 128 Faculty of Science, Okayama University, Okayama, Japan
- 129 Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, OK, USA
- 130 Department of Physics, Oklahoma State University, Stillwater, OK, USA
- 131 Palacký University, RCPTM, Joint Laboratory of Optics, Olomouc, Czech Republic
- 132 Institute for Fundamental Science, University of Oregon, Eugene, OR, USA
- 133 Graduate School of Science, Osaka University, Osaka, Japan
- 134 Department of Physics, University of Oslo, Oslo, Norway
- 135 Department of Physics, Oxford University, Oxford, UK
- 136 LPNHE, Sorbonne Université, Université de Paris, CNRS/IN2P3, Paris, France
- 137 Department of Physics, University of Pennsylvania, Philadelphia, PA, USA
- 138 Konstantinov Nuclear Physics Institute of National Research Centre “Kurchatov Institute”, PNPI, St. Petersburg, Russia
- 139 Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, PA, USA
- 140 (a) Laboratório de Instrumentação e Física Experimental de Partículas - LIP, Lisboa, Portugal; (b) Departamento de Física, Faculdade de Ciências, Universidade de Lisboa, Lisboa, Portugal; (c) Departamento de Física, Universidade de Coimbra, Coimbra, Portugal; (d) Centro de Física Nuclear da Universidade de Lisboa, Lisboa, Portugal; (e) Departamento de Física, Universidade do Minho, Braga, Portugal; (f) Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada, Spain; (g) Dep Física and CEFITEC of Faculdade de Ciências e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal; (h) Instituto Superior Técnico, Universidade de Lisboa, Lisboa, Portugal
- 141 Institute of Physics of the Czech Academy of Sciences, Prague, Czech Republic
- 142 Czech Technical University in Prague, Prague, Czech Republic
- 143 Charles University, Faculty of Mathematics and Physics, Prague, Czech Republic
- 144 Particle Physics Department, Rutherford Appleton Laboratory, Didcot, UK
- 145 IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France
- 146 Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, CA, USA
- 147 (a) Departamento de Física, Pontificia Universidad Católica de Chile, Santiago, Chile; (b) Universidad Andres Bello, Department of Physics, Santiago, Chile; (c) Instituto de Alta Investigación, Universidad de Tarapacá, Arica, Chile; (d) Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile
- 148 Department of Physics, University of Washington, Seattle, WA, USA
- 149 Department of Physics and Astronomy, University of Sheffield, Sheffield, UK
- 150 Department of Physics, Shinshu University, Nagano, Japan
- 151 Department Physik, Universität Siegen, Siegen, Germany
- 152 Department of Physics, Simon Fraser University, Burnaby, BC, Canada
- 153 SLAC National Accelerator Laboratory, Stanford, CA, USA
- 154 Physics Department, Royal Institute of Technology, Stockholm, Sweden

- 155 Departments of Physics and Astronomy, Stony Brook University, Stony Brook, NY, USA
- 156 Department of Physics and Astronomy, University of Sussex, Brighton, UK
- 157 School of Physics, University of Sydney, Sydney, Australia
- 158 Institute of Physics, Academia Sinica, Taipei, Taiwan
- 159 ^(a)E. Andronikashvili Institute of Physics, Iv. Javakishvili Tbilisi State University, Tbilisi, Georgia; ^(b)High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
- 160 Department of Physics, Technion, Israel Institute of Technology, Haifa, Israel
- 161 Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel
- 162 Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece
- 163 International Center for Elementary Particle Physics and Department of Physics, University of Tokyo, Tokyo, Japan
- 164 Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan
- 165 Department of Physics, Tokyo Institute of Technology, Tokyo, Japan
- 166 Tomsk State University, Tomsk, Russia
- 167 Department of Physics, University of Toronto, Toronto, ON, Canada
- 168 ^(a)TRIUMF, Vancouver, BC, Canada; ^(b)Department of Physics and Astronomy, York University, Toronto, ON, Canada
- 169 Division of Physics and Tomonaga Center for the History of the Universe, Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan
- 170 Department of Physics and Astronomy, Tufts University, Medford, MA, USA
- 171 Department of Physics and Astronomy, University of California Irvine, Irvine, CA, USA
- 172 Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden
- 173 Department of Physics, University of Illinois, Urbana, IL, USA
- 174 Instituto de Física Corpuscular (IFIC), Centro Mixto Universidad de Valencia - CSIC, Valencia, Spain
- 175 Department of Physics, University of British Columbia, Vancouver, BC, Canada
- 176 Department of Physics and Astronomy, University of Victoria, Victoria, BC, Canada
- 177 Fakultät für Physik und Astronomie, Julius-Maximilians-Universität Würzburg, Würzburg, Germany
- 178 Department of Physics, University of Warwick, Coventry, UK
- 179 Waseda University, Tokyo, Japan
- 180 Department of Particle Physics, Weizmann Institute of Science, Rehovot, Israel
- 181 Department of Physics, University of Wisconsin, Madison, WI, USA
- 182 Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal, Germany
- 183 Department of Physics, Yale University, New Haven, CT, USA
- 184 Yerevan Physics Institute, Yerevan, Armenia
- ^a Also at Borough of Manhattan Community College, City University of New York, New York NY, USA
- ^b Also at CERN, Geneva, Switzerland
- ^c Also at CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille, France
- ^d Also at Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève, Switzerland
- ^e Also at Departament de Física de la Universitat Autònoma de Barcelona, Barcelona, Spain
- ^f Also at Department of Applied Physics and Astronomy, University of Sharjah, Sharjah, United Arab Emirates,
- ^g Also at Department of Financial and Management Engineering, University of the Aegean, Chios, Greece
- ^h Also at Department of Physics and Astronomy, Michigan State University, East Lansing MI, USA
- ⁱ Also at Department of Physics and Astronomy, University of Louisville, Louisville, KY, USA
- ^j Also at Department of Physics, Ben Gurion University of the Negev, Beer Sheva, Israel
- ^k Also at Department of Physics, California State University, East Bay, USA
- ^l Also at Department of Physics, California State University, Fresno, USA
- ^m Also at Department of Physics, California State University, Sacramento, USA
- ⁿ Also at Department of Physics, King's College London, London, UK
- ^o Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg, Russia
- ^p Also at Department of Physics, Stanford University, Stanford CA, USA
- ^q Also at Department of Physics, University of Adelaide, Adelaide, Australia
- ^r Also at Department of Physics, University of Fribourg, Fribourg, Switzerland
- ^s Also at Department of Physics, University of Michigan, Ann Arbor MI, USA

- ^t Also at Dipartimento di Matematica, Informatica e Fisica, Università di Udine, Udine, Italy
- ^u Also at Faculty of Physics, M.V. Lomonosov Moscow State University, Moscow, Russia
- ^v Also at Giresun University, Faculty of Engineering, Giresun, Turkey
- ^w Also at Graduate School of Science, Osaka University, Osaka, Japan
- ^x Also at Hellenic Open University, Patras, Greece
- ^y Also at IJCLab, Université Paris-Saclay, CNRS/IN2P3, 91405, Orsay, France
- ^z Also at Institucio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain
- ^{aa} Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany
- ^{ab} Also at Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands
- ^{ac} Also at Institute for Nuclear Research and Nuclear Energy (INRNE) of the Bulgarian Academy of Sciences, Sofia, Bulgaria
- ^{ad} Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary
- ^{ae} Also at Institute of Particle Physics (IPP), Vancouver, Canada
- ^{af} Also at Institute of Physics, Academia Sinica, Taipei, Taiwan
- ^{ag} Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan
- ^{ah} Also at Institute of Theoretical Physics, Ilia State University, Tbilisi, Georgia
- ^{ai} Also at Instituto de Fisica Teorica, IFT-UAM/CSIC, Madrid, Spain
- ^{aj} Also at Joint Institute for Nuclear Research, Dubna, Russia
- ^{ak} Also at Louisiana Tech University, Ruston LA, USA
- ^{al} Also at LPNHE, Sorbonne Université, Université de Paris, CNRS/IN2P3, Paris, France
- ^{am} Also at Manhattan College, New York NY, USA
- ^{an} Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia
- ^{ao} Also at National Research Nuclear University MEPhI, Moscow, Russia
- ^{ap} Also at Physics Department, An-Najah National University, Nablus, Palestine
- ^{aq} Also at Physics Dept, University of South Africa, Pretoria, South Africa
- ^{ar} Also at Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg, Germany
- ^{as} Also at School of Physics, Sun Yat-sen University, Guangzhou, China
- ^{at} Also at The City College of New York, New York NY, USA
- ^{au} Also at TRIUMF, Vancouver BC, Canada
- ^{av} Also at Università di Napoli Parthenope, Napoli, Italy.
- Deceased: O. Igonkina, B.T. King, D. Lellouch.